

Kap 9+10 Rotasjon, spinn.

Eksempler og
demonstrasjoner i forelesning

(N1-rot):

Bevaring av spinn i kontorstol



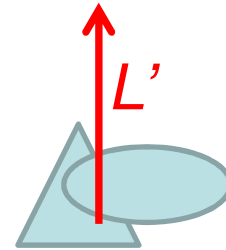
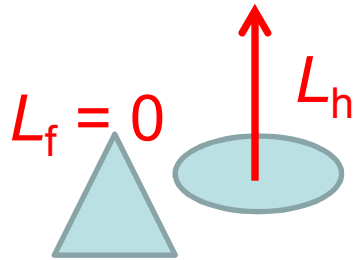
= stol+ foreleser, spinn L_f



= hjul, spinn L_h

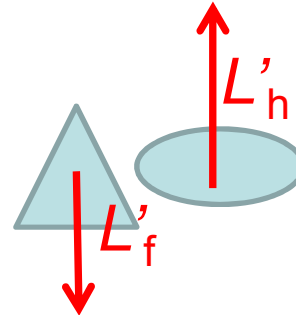
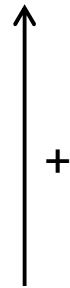
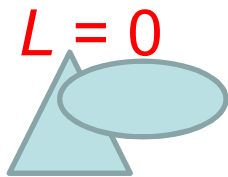
$$L = L_f + L_h$$

1



$$L' = L_h$$

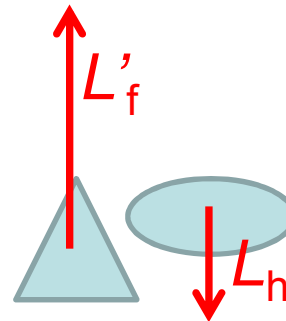
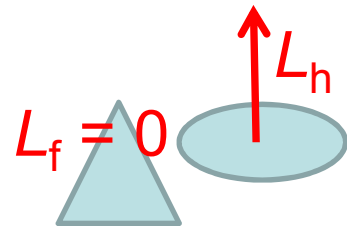
2



$$0 = L'_h - L'_f$$

$$\Rightarrow L'_f = L'_h$$

3



$$L_h = L'_f - L_h$$

$$\Rightarrow L'_f = 2L_h$$

Spinn: $L = I \omega = \text{konstant!}$

Personer inn mot sentrum:

$$I = \sum m_i r_i^2$$

avtar

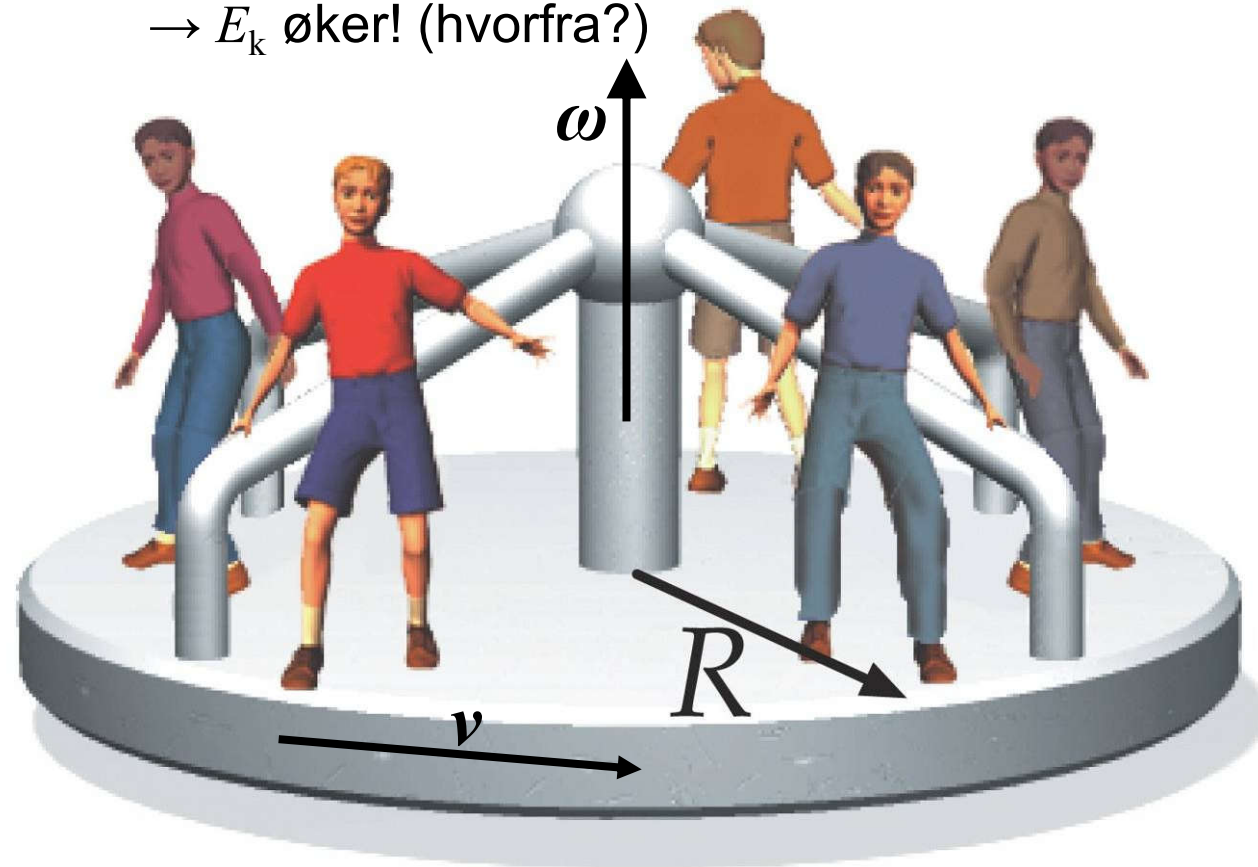
ω
må øke!

**Ikke stivt
legeme!**

Kinetisk energi: $E_k = \frac{1}{2} I \omega^2 = \frac{1}{2} L \omega$

→ L konstant, ω øker

→ E_k øker! (hvorfra?)



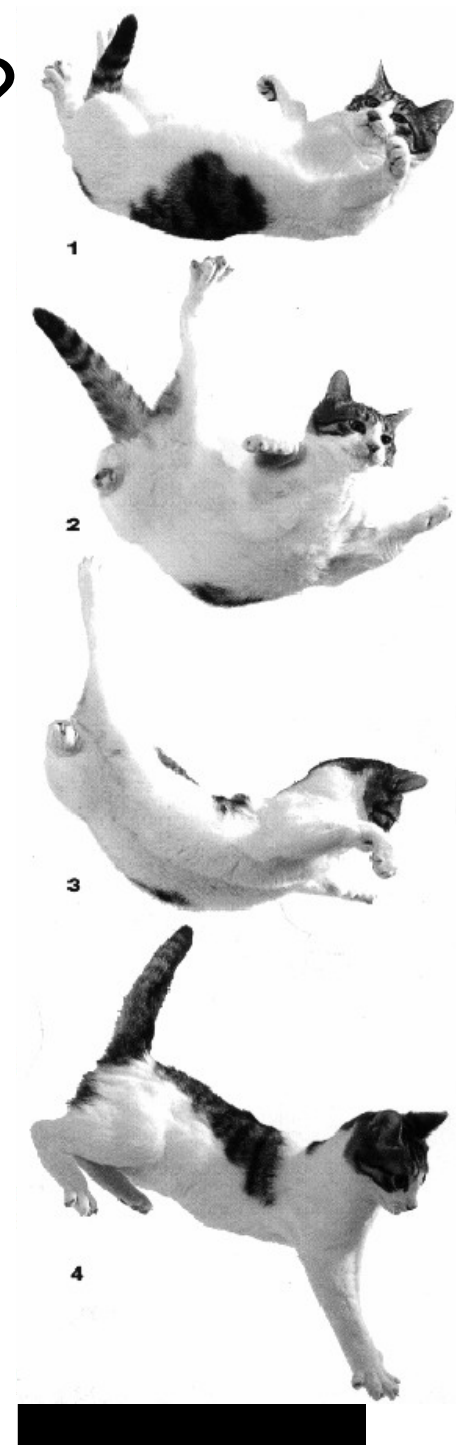
Spinn for fallende katt bevart?

Katter lander

- alltid på føttene!

$L = 0$ ved start og ved slutt

$L = 0$ underveis !?

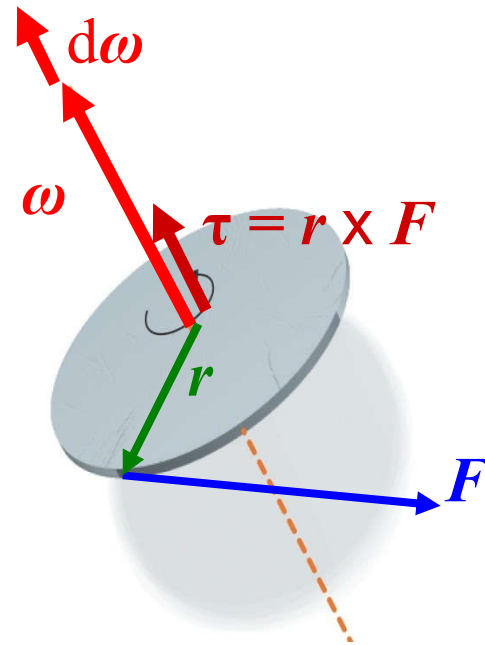


Gyroskop

1. Lodd holder hjulet i balanse
2. $L = I\omega$ konstant (uten τ_{ytre})
→ gyrokompass
3. Stor motstand mot endring
4. Endring av akseretning ved kraft normalt på endringen



(N2-rot)



Alle legemer:

$$\boldsymbol{\tau} = d\mathbf{L} / dt \quad (10.29)$$

Stivt legeme
om symmetriakse:

$$\mathbf{L} = I \boldsymbol{\omega}$$

=>

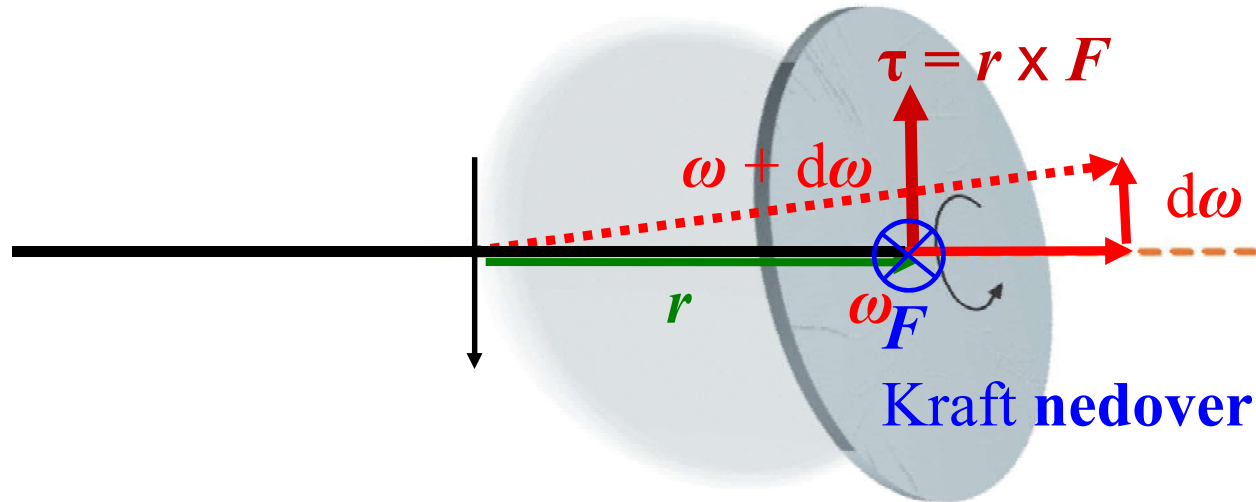
$$\boldsymbol{\tau} = I d\boldsymbol{\omega} / dt \quad (10.7)$$

Raskere rotasjon om samme akse:
 $\boldsymbol{\omega} \rightarrow \boldsymbol{\omega} + d\boldsymbol{\omega}$ alle i samme retning
(N2-rot): $\boldsymbol{\tau} dt = I d\boldsymbol{\omega}$
=> $\boldsymbol{\tau}$ i samme retning som $d\boldsymbol{\omega}$
=> \mathbf{F} som i figuren

Hva hvis akseretningen skal endres?

Endring av akseretning

Sett ovenfra:



Alle legemer:

$$\tau = dL / dt \quad (10.29)$$

Stivt legeme
om symmetriakse:

$$L = I \omega$$

\Rightarrow

$$\tau = I d\omega / dt \quad (10.7)$$

Endring akseretning:

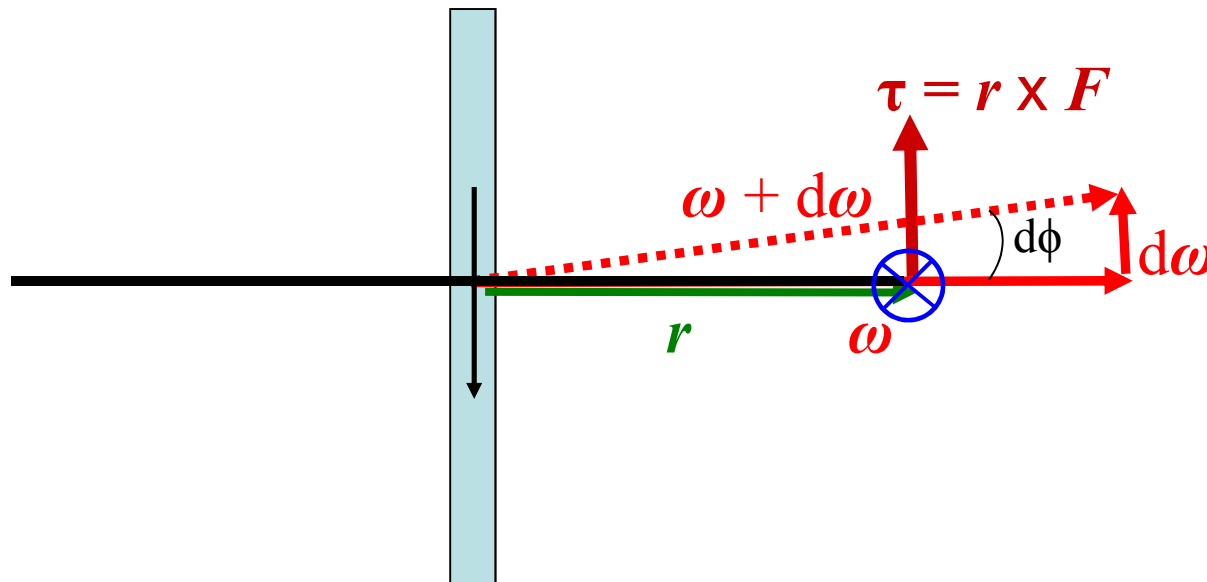
$$\omega \rightarrow \omega + d\omega$$

(N2-rot): $\tau dt = I d\omega$

$\Rightarrow \tau$ i samme retning som $d\omega$

$\Rightarrow F$ nedover

Med vedvarende F får vi
presesjon



Alle legemer:

$$\tau = dL / dt \quad (10.29)$$

Stivt legeme
om symmetriakse:

$$L = I \omega$$

\Rightarrow

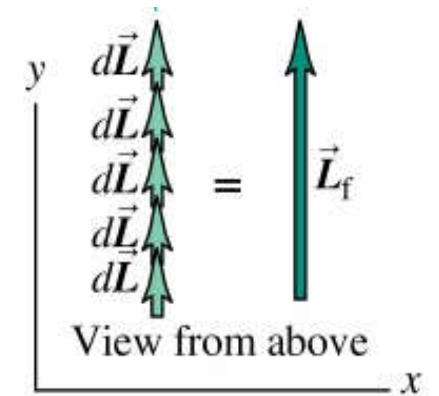
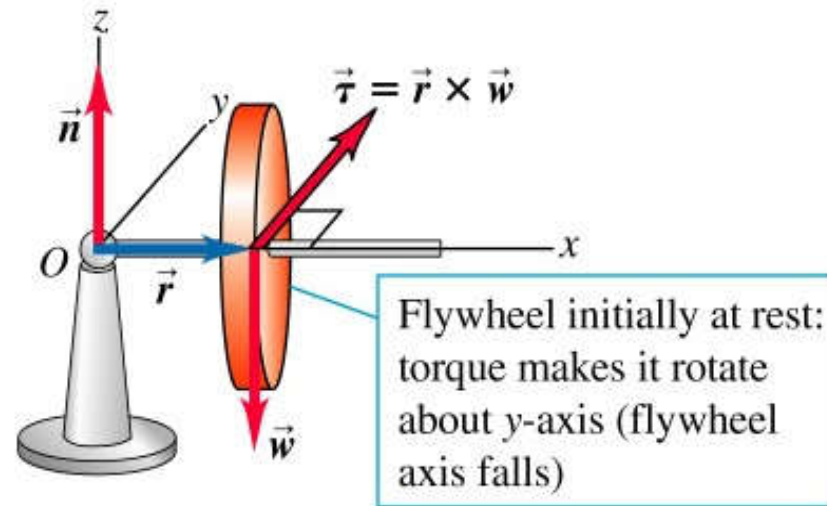
$$\tau = I d\omega / dt \quad (10.7)$$

$$d\phi = \frac{d\omega}{\omega}$$

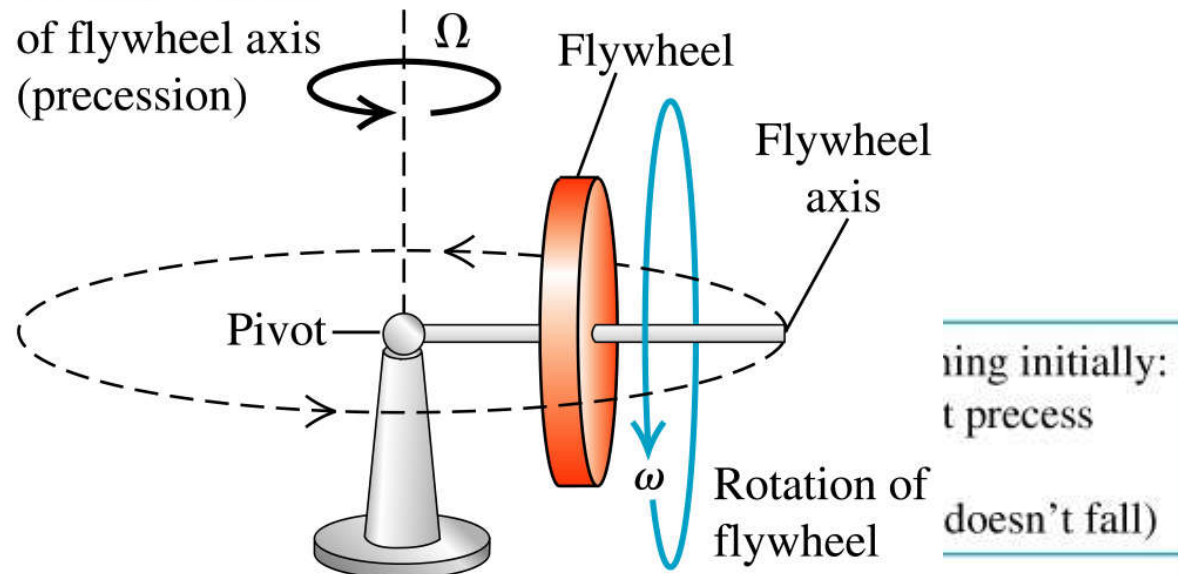
Presesjonsfrekvens: $\Omega_p = \frac{d\phi}{dt} = \frac{d\omega}{dt} \frac{1}{\omega} \stackrel{(N2-rot)}{=} \frac{\tau}{I} \frac{1}{\omega} = \frac{Fr}{I\omega}$

Sykkelhjul

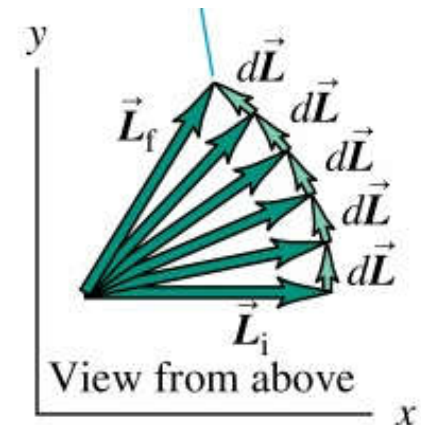
Ikke-roterende hjul:



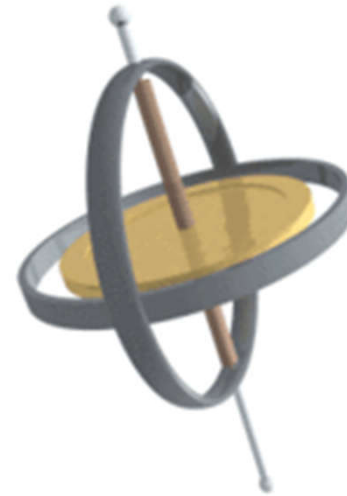
Circular motion of flywheel axis (precession)



Sett ovenfra:



Snurrebasser



Alle legemer:

$$\boldsymbol{\tau} = d\mathbf{L} / dt \quad (10.29)$$

Stivt legeme
om symmetriakse:

$$\mathbf{L} = I \boldsymbol{\omega}$$

\Rightarrow

$$\boldsymbol{\tau} = I d\boldsymbol{\omega} / dt \quad (10.7)$$

Rotasjon av stive legemer

- Tregghetsmoment $I = \sum r_i^2 m_i$ (om en gitt akse)
- Rotasjonsenergi $E_k = \frac{1}{2} \sum m_i v_i^2 = \frac{1}{2} I \omega^2$
- Kraftmoment: $\boldsymbol{\tau} = \mathbf{r} \times \mathbf{F}$
- Spinn (dreieimpuls) $\mathbf{L} = \mathbf{r} \times m \mathbf{v} = I \boldsymbol{\omega}$
- Spinnsatsen (N2-rot): $\boldsymbol{\tau} = d\mathbf{L} / dt = I d\boldsymbol{\omega} / dt$
- Ingen ytre moment (N1-rot): $\mathbf{L} = \text{konst.}$

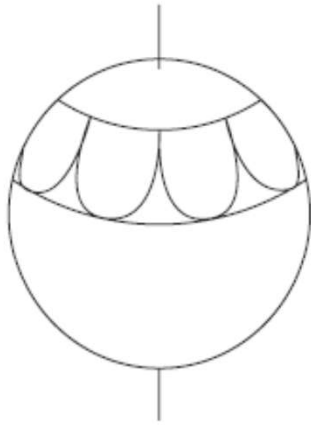
stivt legeme om
sym.akse:

Matematisk forklaring av fysikken ofte eneste mulige

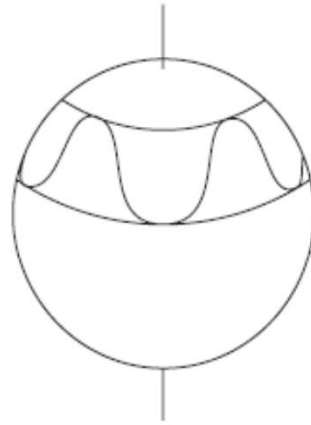
Richard Feynman (am. fysiker/pedagog, 1918-1988):

”...many simple things can be deduced mathematically more rapidly than they can really be understood in a fundamental or simple sense. This is a strange characteristic, and as we get into more and more advanced work there are circumstances in which mathematics will produce results which *no one* has really been able to understand in any direct fashion.”

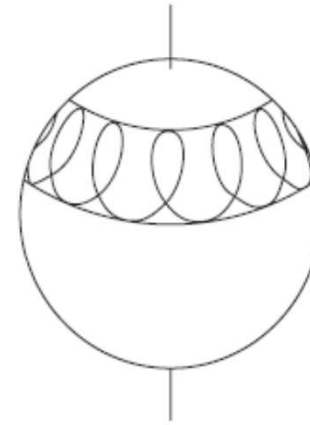
Nutasjon



(A) Released from rest



(B) Released with forward speed



(C) Released with backward speed

Hva betyr gyroeffekten for å holde sykkel oppe?

Mr. Jones testet dette med hjul som roterte motsatt retning, dvs. motsatt gyroeffekt.

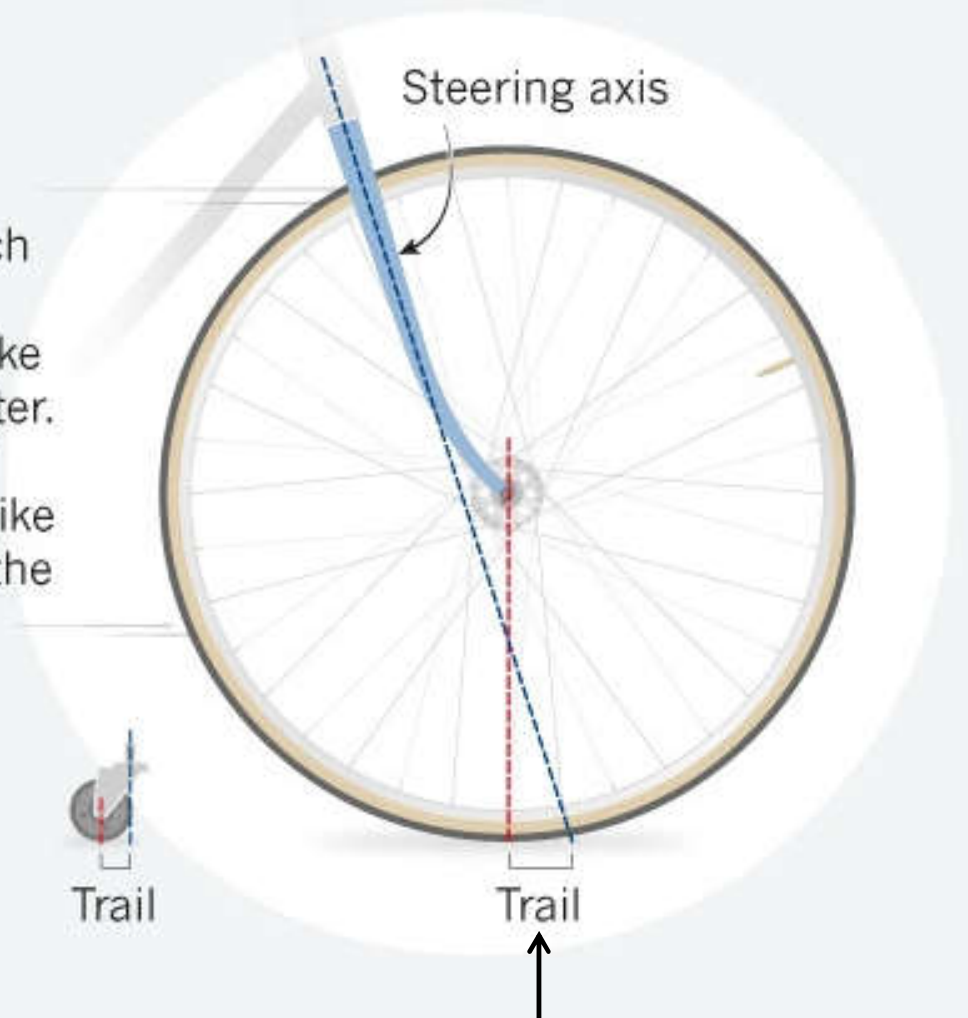
=> En URB (UnRidableBicycle)?



Gyroeffekten
ingen
praktisk
betydning!

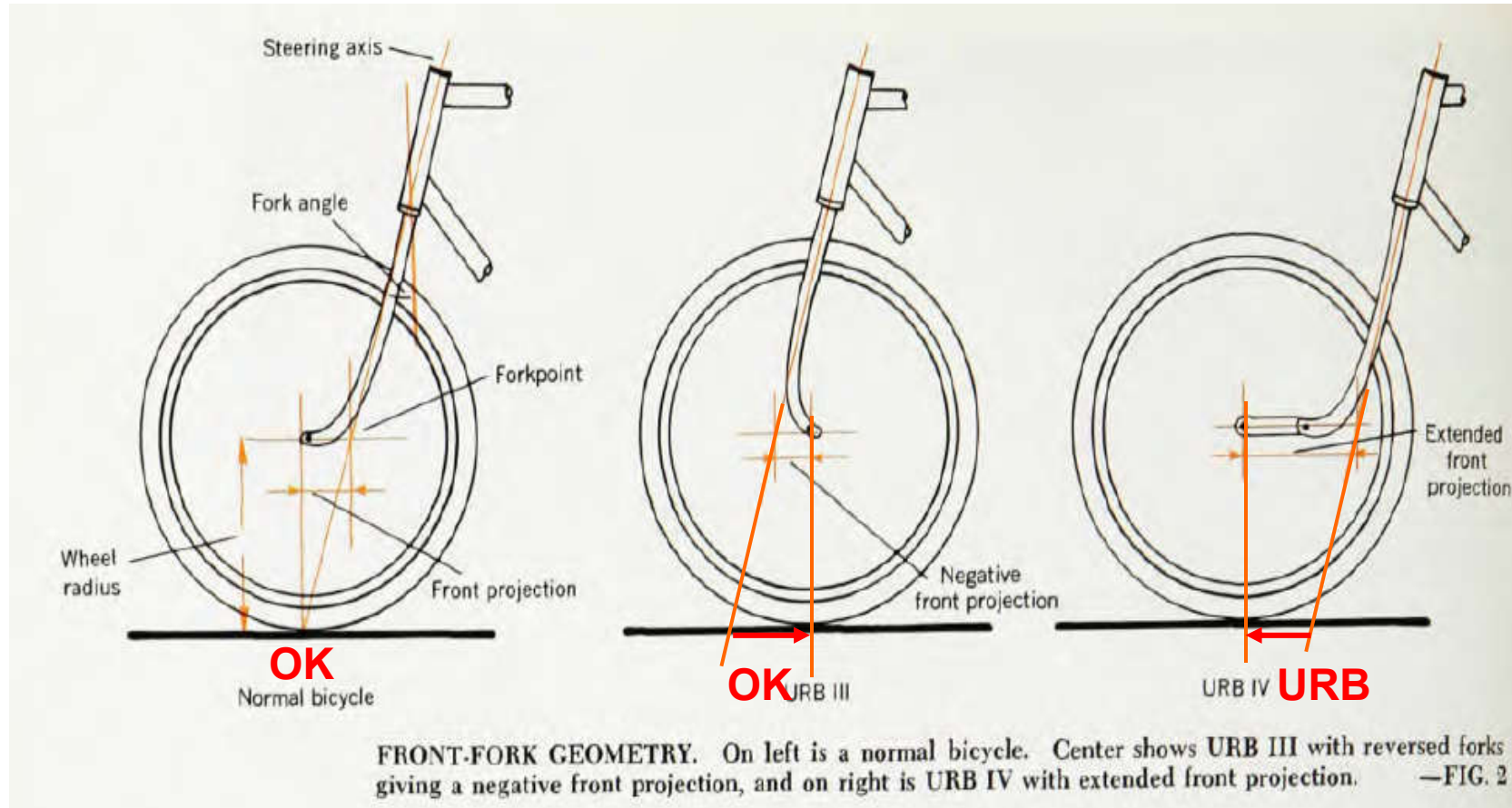
THE CASTER TRAIL

A bicycle's front-wheel steering axis sits slightly ahead of the point at which the wheel touches the ground, creating a 'trail' like that of an office-chair caster. This means the wheel will turn in the direction the bike is travelling (or falling, as the case may be).



Viktig at bakkekontakt er bak styreaksen

Mr. Jones' URB-sykler:



Snøsykler (snowbikes): Null gyroeffekt



Fra wikipedia

For spesielt interesserte:

Sykkelens stabilitet, referanser.

Noen av lenkene krever IP-adresse fra NTNU for å få tilgang.

Kooijman et al: A Bicycle Can Be Self-Stable Without Gyroscopic or Caster Effects
Science 332 (2011), pp. 339-342 (med mange flere referanser)

<http://science.sciencemag.org/content/332/6027/339.full>

B. Borrell: The bicycle problem that nearly broke mathematics,
Nature 535, 21.July 2016, pp. 338-342.

<http://www.nature.com/news/the-bicycle-problem-that-nearly-broke-mathematics-1.20281>

J. Matson: A Bicycle Built for None: What Makes a Riderless Bike Stable?
Scientific American April 2011

<http://www.scientificamerican.com/article/self-stable-bike/>

D.E.H. Jones: The stability of the bicycle,
Physics Today, April 1970, pp. 34-40

http://www.phys.lsu.edu/faculty/gonzalez/Teaching/Phys7221/vol59no9p51_56.pdf
<http://home.phys.ntnu.no/brukdef/undervisning/tfy4145/arkiv/2010/diverse/UnridableBicycle.pdf>

J. Lowell and H.D. McKell, The Stability of Bicycles, (teknisk avansert)
Am. J. Phys. 50 (1982), pp. 1106-1112.

https://www.me.utexas.edu/~longoria/VSDC/09_vehicle_roll_dynamics_and_control/09_3b_Lowell_1982_The%20Stability%20of%20Bicycles.pdf

Bicycle wheel / Turning a bicycle:

<http://hyperphysics.phy-astr.gsu.edu/hbase/mechanics/bicycle.html#c2>