

fact dik em py kanda milith.

1

I

$$\lambda = \frac{h}{mv} = \left(\frac{h^2}{2meV} \right)^{1/2}$$

y CGS : (1 volt = $\frac{1}{300}$ CGS potential unit)

$$\lambda = \left[\frac{6.623 \cdot 10^{-12}}{2 \cdot 9.106 \cdot 10^{-18} \cdot 1.6 \cdot 10^{-10}} \cdot \frac{300}{V(\text{volt})} \right]^{1/2} \text{ centimeter}$$

$$= \left[\dots \right] \cdot 10^8 \text{ } \overset{\circ}{\text{A}}$$

$$= \left[\frac{43,8641 \cdot 10^{-54} \cdot 300 \cdot 10^{16}}{87,434 \cdot 10^{-38} V(\text{volt})} \right]^{1/2} = \sqrt{\frac{150}{V(\text{volt})}}$$

$$\therefore \text{ knt} = 150$$

i SI - unit

$$\lambda = \left[\frac{(6.623 \cdot 10^{-34})^2}{2 \cdot 9.106 \cdot 10^{-31} \cdot 1.6 \cdot 10^{-19} V} \right]^{1/2} \text{ meter}$$

$$= \left[\frac{43,8641 \cdot 10^{-68}}{29,372 \cdot 10^{-50} V} \cdot 10^{20} \right]^{1/2} \text{ Angstrom}$$

$$= \sqrt{\frac{150}{V(\text{volt})}}$$

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Facit II :

Impulsstörzfunktion

$$a) \phi(k, 0) = \frac{1}{\sqrt{2\pi}} \int \psi(x, 0) dx \cdot e^{-\frac{i}{\hbar} k x} = \left(\frac{1}{\sqrt{2\pi}} \right)^{1/2} \left(\frac{1}{\sqrt{2\pi}} \right)^{1/2} e^{-\frac{i}{\hbar} k x}$$

$$\omega(k, 0) = |\phi(k)|^2 = \left(\frac{1}{\sqrt{2\pi}} \right)^{-1/2} \cdot \frac{1}{\sqrt{2\pi}} \cdot e^{-\frac{i}{\hbar} k x} = 2$$

$$\therefore \underline{\Delta p^2 = \frac{1}{4} \hbar^2 / \Delta x^2} \quad \omega(p) = \left(\frac{1}{\sqrt{2\pi}} \right)^{-1/2} e^{-\frac{p^2}{2\Delta p^2}}$$

$$b) \bar{\psi}(x, t) = \frac{1}{\sqrt{2\pi}} \int \phi(k, 0) \cdot e^{-\frac{i\hbar k^2 t}{2m}} dp \cdot e^{\frac{i\hbar k x}{\hbar}}$$

$$\int e^{-k^2 \left(\frac{\Delta x^2}{\hbar^2} + \frac{i\hbar t}{2m} \right) + \frac{i\hbar k x}{\hbar}} = \left(\frac{\hbar^2 t^2}{\Delta x^2 + \frac{i\hbar t}{2m}} \right)^{1/2} \cdot e^{-\frac{x^2}{2 \left(\Delta x^2 + \frac{i\hbar t}{2m} \right)}}$$

$$\frac{1}{\Delta x^2 + \frac{i\hbar t}{2m}} = \frac{\Delta x^2 - \frac{i\hbar t}{2m}}{\Delta x^4 + \frac{\hbar^2 t^2}{4m^2}} = \frac{1 - \frac{i\hbar t}{2m \Delta x^2}}{\Delta x^2 + \frac{\Delta p^2 \cdot t^2}{m^2}} = \frac{1 - \frac{i\hbar t}{2m \Delta x^2}}{\Delta x^2(t)}$$

$$\Rightarrow \omega(x, t) = \left[\frac{1}{\Delta x^2(t)} \right]^{-1/2} \cdot e^{-\frac{x^2}{2 \Delta x^2(t)}}$$

hier $\Delta x^2(t) = \Delta x^2 + \left(\frac{\hbar t}{m} \right)^2 \cdot t^2$

$$c) \psi(x, 0) = \left(\frac{1}{\sqrt{2\pi}} \right)^{1/2} \cdot e^{-\frac{(x-x_0)^2}{4\Delta x^2} + \frac{i}{\hbar} p_0 x}$$

for example.

(allmähliche $\psi = f(x-x_0) \cdot e^{\frac{i}{\hbar} p_0 x}$)

und f symmetrisch: $|f|^2(z) = |f|^2(-z)$

d) kann man mit dem für man in dem Fall ψ^0 für die Galilei Transformation $x = x_0 + \frac{p_0}{m} t$ oder a) b) \Rightarrow

$$\omega(p) = N e^{-\frac{(p-p_0)^2}{2\Delta p^2}}, \quad \omega(x, t) = N' e^{-\frac{(x-x_0 - \frac{p_0}{m} t)^2}{2\Delta x^2(t)}}$$

$$e) \frac{d}{dt} \langle p^2 \rangle = \left\langle \frac{i}{\hbar} [H, p^2] \right\rangle = 0 \quad \text{für } H = \frac{p^2}{2m}$$

$$f) \psi = \rho e^{i\delta(x)}, \quad \frac{\hbar}{i} \nabla \psi = \frac{\hbar}{i} (\nabla \rho + i\rho \nabla \delta) \cdot e^{i\delta}$$

$$\langle \sigma \rangle = \frac{1}{m} \frac{\hbar}{i} \left\{ \int \rho \nabla \rho + i \int \rho^2 \nabla \delta \right\}$$

früher $\langle \sigma \rangle = 0$ das ist für Bewegung mit v_0 in δ und δ in σ das heißt δ ist konstant.

$$\begin{aligned}
 \text{b)} \quad \psi(x; t + \Delta t) &= \int K(x, s; \Delta t) \psi(s, t) \\
 &= e^{-\frac{i}{\hbar} V(x) \Delta t} \int K^0 ds \left[\sum_{n=0}^{\infty} \frac{(s-x)^n}{n!} \frac{\partial^n \psi}{\partial s^n} \Big|_{s=x} \right] \\
 &= \psi(x, t) + \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} \Delta t + \mathcal{O}(\Delta t)^2 \\
 &= \psi(x, t) + \frac{i}{\hbar} \left(\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} - V\psi \right) \Delta t + \mathcal{O}(\Delta t)^2
 \end{aligned}$$

v. s. $\psi(x, t) + \Delta t \cdot \frac{\partial \psi}{\partial t} + \mathcal{O}(\Delta t)^2 = \text{h. s.}$

$$\begin{aligned}
 \psi(x, t + \Delta t) - \psi(x, t) &= \\
 \Delta t \frac{\partial \psi}{\partial t} + \mathcal{O}(\Delta t)^2 &= -\frac{i}{\hbar} \left(\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi \right) \Delta t + \mathcal{O}(\Delta t)^2
 \end{aligned}$$

denn $i\hbar \frac{\partial \psi}{\partial t} = \mathcal{H}\psi$

$$\begin{aligned}
 \text{a)} \quad K^0: \text{Zur } (i\hbar \frac{\partial}{\partial t} - H^0) K^0 &= 0 \quad \text{für } t > 0 & (1) \\
 K^0 &= 0 \quad t < 0 & (2) \\
 \text{oder } K^0 &\rightarrow \delta(t-t_0) \text{ mit } t \rightarrow t_0 & (3)
 \end{aligned}$$

Zur: (1) $(i\hbar \frac{\partial}{\partial t} - H^0) K^0 = \delta(t-t_0) \delta(x-x_0)$

(2') $K^0 = 0$ für $t < t_0$

$$K^0 = \int a(k, t) dk e^{i k(R - \omega T)} \quad \text{mit } R = x - x_0 \quad T = t - t_0$$

(1') $\Rightarrow \int [\omega - \frac{\hbar^2 k^2}{2m}] a(k, \omega) \exp. \Rightarrow a = \frac{1}{(2\pi)} + \frac{1}{\omega - \frac{\hbar^2 k^2}{2m}}$

(2') $\Rightarrow K = \lim_{\epsilon \rightarrow 0} \frac{1}{(2\pi)} \int dk e^{i k R} e^{-i \frac{\hbar^2 k^2}{2m} T} \int_{-\infty}^{\infty} \frac{d\omega e^{-i\omega T}}{\omega + i\epsilon - \frac{\hbar^2 k^2}{2m}}$

Das ist für $T > 0$ anders 0.

$$\int e^{i k R} d\Omega(\hat{k}) = 4\pi \frac{m^2 \hbar^2}{k^2} \int d^3k \exp = \frac{4\pi}{k^2} \int_0^{\infty} k^2 dk \sin k R e^{-\frac{i\hbar^2 k^2}{2m} T}$$

$$-\frac{4\pi}{k^2} \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i k R} \dots = \left(\frac{2m}{\hbar^2} \right)^{3/2} e^{i \frac{m R^2}{2\hbar T}} \quad \text{g. e. d.}$$