

FY8201 / TFY8 Nanoparticle and polymer physics I

SOLUTION of EXERCISE 2

Eq. (x.x) refers to version AM24nov05 of lecture notes: “Nanoparticle and polymer physics”.

First we repeat given equations:

$$L^{(p)} := \lim_{N \rightarrow \infty} \left\langle |Q_1|^{-1} \vec{Q}_1 \cdot \sum_{k=1}^{N-1} \vec{Q}_k \right\rangle = \lim_{N \rightarrow \infty} |Q_1|^{-1} \sum_{k=1}^{N-1} \left\langle \vec{Q}_1 \cdot \vec{Q}_k \right\rangle \quad (1)$$

$$L^{(p)} = Q \lim_{N \rightarrow \infty} \sum_{k=1}^{N-1} [1, 0, 0] \left\langle \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \right\rangle \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \quad (2)$$

$$\vec{Q}_k = \vec{\Omega}_1^{(\xi)} \cdot \vec{\Omega}_2^{(\xi)} \cdots \vec{\Omega}_{k-1}^{(\xi)} \cdot Q \cdot \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = Q \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \cdot \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \quad (3)$$

$$L^{(p)} = Q \lim_{N \rightarrow \infty} \sum_{k=1}^{N-1} [1, 0, 0] \prod_{m=1}^{k-1} \left\langle \vec{\Omega}_m^{(\xi)} \right\rangle \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}. \quad (4)$$

We assume low concentrations of polymers so that only intramolecular (and no intermolecular) forces determine the molecular configuration \vec{Q}_k . Further we assume that there are interaction only between the nearest neighbour segments in the polymer chain. Thermal average is then expressed by

$$\langle B \rangle = \frac{\int \cdots \int B \exp\{-\mathcal{H}(\vec{P}, \vec{Q})\beta\} d\vec{P} d\vec{Q}}{\int \cdots \int \exp\{-\mathcal{H}(\vec{P}, \vec{Q})\beta\} d\vec{P} d\vec{Q}} \quad (5)$$

where $\beta = (k_B T)^{-1}$ and the Hamiltonian \mathcal{H} is a function of nearest neighbour vectors

$$\mathcal{H} = \sum_{i=1}^{N-2} h_i(\vec{Q}_i, \vec{Q}_{i+1}) = \sum_{i=1}^{N-2} h_i(\vec{Q}_i, \vec{\Omega}_i^{(\xi)} \cdot \vec{Q}_{i+1}) \quad (6)$$

where h_i are any functions (not required to be specified). Note the notation: In the first sum \vec{Q}_i and \vec{Q}_{i+1} are given relative to their respective local coordinate systems, while in the last sum they are given relative to the *same* coordinate system, namely the coordinate system of segment vector number i .

The thermal averaging in Eq. (2) now yields

$$\left\langle \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \right\rangle = \frac{1}{Z} \int \int \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \exp\{-\mathcal{H}\beta\} d\vec{P} d\vec{Q} \quad (7)$$

where Z is the normalization constant (the partition function). Insertion of \mathcal{H} from Eq. (6) yields

$$\begin{aligned} \left\langle \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \right\rangle &= \frac{1}{Z} \int \int \left(\prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \right) \exp\left\{-\beta \sum_{i=1}^{N-2} h_i(\vec{Q}_i, \vec{\Omega}_i^{(\xi)} \cdot \vec{Q}_{i+1})\right\} d\vec{P} d\vec{Q} \\ &= \frac{1}{Z'} \int \int \prod_{m=1}^{k-1} \left(\vec{\Omega}_m^{(\xi)} \exp\{-\beta \cdot h_m(\vec{Q}_m, \vec{\Omega}_m^{(\xi)} \cdot \vec{Q}_{m+1})\} \right) d\vec{P} d\vec{Q} \end{aligned} \quad (8)$$

where all integrals containing h_i with $i \in [k, N-2]$ have been cancelled by the corresponding integrals in Z . Noting that only h_1 contains $\vec{\Omega}_1^{(\xi)}$, h_2 contains $\vec{\Omega}_2^{(\xi)}$ etc., we obtain:

$$\left\langle \prod_{m=1}^{k-1} \vec{\Omega}_m^{(\xi)} \right\rangle = \frac{1}{Z'} \prod_{m=1}^{k-1} \left(\int \int \vec{\Omega}_m^{(\xi)} \exp\{-\beta \cdot h_m(\vec{Q}_m, \vec{\Omega}_m^{(\xi)} \vec{Q}_{m+1})\} d\vec{P} d\vec{Q} \right) \quad (9)$$

$$= \prod_{m=1}^{k-1} \left\langle \vec{\Omega}_m^{(\xi)} \right\rangle. \quad (10)$$

B) In the Kirkwood-Riseman chain \vec{Q}_{k+1} can be rotated freely an angle ξ_{k2} around \vec{Q}_k with a fixed angle $\xi_{k1} = \xi$. As the angle ξ_{k2} is free to rotate between 0 and 2π , the Hamiltonian \mathcal{H} does not depend on ξ_{k2} and the average is zero:

$$\langle \cos \xi_{k2} \rangle = \frac{1}{Z} \int \cdots \int \cos \xi_{k2} \exp\{-\mathcal{H}(\vec{P}, \vec{Q})\beta\} d\xi_{k2} d\vec{P} d\vec{Q} = 0 \quad (11)$$

$$(12)$$

And similarly $\langle \sin \xi_{k2} \rangle = 0$. Therefore (from Eq. (2.53))

$$\left\langle \vec{\Omega}_m^{(\xi)} \right\rangle = \begin{bmatrix} \cos \xi & -\sin \xi & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \equiv \vec{\Omega} \quad (13)$$

From the result in A) of this exercise we get

$$\left\langle \prod_{m=n}^p \vec{\Omega}_m^{(\xi)} \right\rangle = \prod_{m=n}^p \left\langle \vec{\Omega}_m^{(\xi)} \right\rangle = \prod_{m=n}^p \vec{\Omega} = \vec{\Omega}^{(p-n+1)} \quad (14)$$

Inserted in Eq. (4) for persistence length:

$$\begin{aligned} L^{(p)} &= Q \lim_{N \rightarrow \infty} \sum_{k=1}^{N-1} [1, 0, 0] \vec{\Omega}^{(k-1)} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = Q \lim_{N \rightarrow \infty} \sum_{k=1}^{N-1} (\cos \xi)^{k-1} \\ &= Q \sum_{k=1}^{\infty} (\cos \xi)^{k-1} = Q(1 + \cos \xi + \cos^2 \xi + \cdots) = \frac{Q}{1 - \cos \xi}. \end{aligned} \quad (15)$$

The sum converges for $|\cos \xi| < 1$.

An alternative calculation according to Doi and Edwards

The average of the projection of \vec{Q}_k on \vec{Q}_{k-1} equals $\langle \vec{Q}_k \cdot \vec{Q}_{k-1} \rangle = Q^2 \cos \xi$, provided all Q_m where $m \neq k$ are being kept constant. Repeating this projection $k-1$ times we obtain $\langle \vec{Q}_k \cdot \vec{Q}_1 \rangle = Q^2 (\cos \xi)^{k-1}$.

From the definition in Eq. (1) we get

$$\begin{aligned} L^{(p)} &= \lim_{N \rightarrow \infty} Q^{-1} \sum_{k=1}^{N-1} \langle \vec{Q}_1 \cdot \vec{Q}_k \rangle \\ &= Q \sum_{k=1}^{\infty} (\cos \xi)^{k-1} = \frac{Q}{1 - \cos \xi}, \end{aligned} \quad (16)$$

as above.