

# Formelliste

## TFY4155/FY1003 Elektrisitet og magnetisme

Formlenes gyldighetsområde og de ulike symbolenes betydning antas å være kjent. Symbolbruk som i forelesningene.  
Siste revisjon: 20.1.12; 5.3.12(retta feil i  $R = \dots$ ) (2 sider).

$Q, \rho$  og  $\sigma$  uten indeks viser til *frie* ladninger.  $Q_i, \rho_i$  og  $\sigma_i$  er indusert ladning.

$I$  og  $\vec{J}$  uten indeks er ledningsstrøm (conducting current),  $I_d$  og  $\vec{J}_d$  er forskyvningsstrøm (displacement current).

$$\text{Coulombs lov: } \vec{F}_{12} = \frac{1}{4\pi\epsilon} \frac{q_1 q_2}{r^2} \hat{\mathbf{r}}_{12} \quad \vec{E} = \frac{1}{4\pi\epsilon} \frac{q}{r^2} \hat{\mathbf{r}}$$

$$\text{Gauss' lov integralform: } \oint \vec{D} \cdot d\vec{A} = Q \quad \oint \vec{E} \cdot d\vec{A} = Q/\epsilon \quad \oint \vec{P} \cdot d\vec{A} = -Q_i \quad \oint \vec{B} \cdot d\vec{A} = 0$$

$$\text{Gauss' lov differensialform: } \operatorname{div} \vec{D} = \rho \quad \operatorname{div} \vec{E} = \rho/\epsilon \quad \operatorname{div} \vec{P} = -\rho_i \quad \operatorname{div} \vec{B} = 0$$

$$\text{Fluks: } \Phi_E = \iint \vec{E} \cdot d\vec{A} \quad \Phi = \iint \vec{D} \cdot d\vec{A} = \epsilon \Phi_E \quad \Phi_B = \iint \vec{B} \cdot d\vec{A}$$

$$\text{Amperes lov: } \oint \vec{B} \cdot d\vec{s} = \mu \left( I + \epsilon \frac{\partial \Phi_E}{\partial t} \right) \quad \oint \vec{H} \cdot d\vec{s} = I + \frac{\partial \Phi}{\partial t} \quad \operatorname{curl} \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\text{Faradays lov: } \mathcal{E} = -\frac{\partial \Phi_B}{\partial t} = -L \frac{dI}{dt} \quad \oint \vec{E} \cdot d\vec{s} = -\frac{\partial \Phi_B}{\partial t} \quad \operatorname{curl} \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\text{Maxwells likninger: } \operatorname{div} \vec{D} = \rho \quad \operatorname{div} \vec{B} = 0 \quad \operatorname{curl} \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \operatorname{curl} \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\text{Forskyvningsstrøm: } I_d = \frac{\partial \Phi}{\partial t}, \quad \vec{J}_d = \frac{\partial \vec{D}}{\partial t}$$

$$\text{Elektrisk dipolmoment: } \vec{p} = q\vec{d} \quad (\text{fra - til +}) \quad \text{Polarisering: } \vec{P} = \frac{\sum \vec{p}}{\text{volum}}$$

$$\text{Magnetisk (dipol)moment: } \vec{\mu} = \vec{m} = IA \quad \text{Magnetisering: } \vec{M} = \frac{\sum \vec{\mu}}{\text{volum}}$$

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon \vec{E} = \epsilon_r \epsilon_0 \vec{E} \quad \vec{P} = \chi_e \epsilon_0 \vec{E} \quad \epsilon_r = 1 + \chi_e$$

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M} = \mu \vec{H} = \mu_r \mu_0 \vec{H} \quad \vec{M} = \chi_m \vec{H} \quad \mu_r = 1 + \chi_m$$

$$\text{Elektrisk potensial: } V_a - V_b = - \int_b^a \vec{E} \cdot d\vec{s}, \quad \vec{E} = -\vec{\nabla}V, \quad \text{Relativt } \infty: \quad V(r) = \int \frac{dq}{4\pi\epsilon r}$$

$$\text{Energi og energitethet: } U = \frac{1}{2} \iiint V dq \quad \text{Elektrisk: } u = \frac{1}{2} \vec{D} \cdot \vec{E} \quad \text{Magnetisk: } u = \frac{1}{2} \vec{B} \cdot \vec{H}$$

$$\text{Kondensatorer: } C = \frac{Q}{V} \quad \text{Kulekondensator: } C = 4\pi\epsilon_0 R \quad \text{Energi: } U = \frac{1}{2} QV = \frac{1}{2} CV^2$$

$$\text{Platekondensator: } C = \epsilon \frac{A}{d} \quad \text{Parallelkopling: } C = \sum_i C_i \quad \text{Seriekopling: } \frac{1}{C} = \sum_i \frac{1}{C_i}$$

$$\text{Kraft på strømførende leder: } d\vec{F} = Id\vec{s} \times \vec{B} \quad \text{Lorentzkrafta: } \vec{F} = q \left( \vec{E} + \vec{v} \times \vec{B} \right)$$

$$\text{Biot-Savarts lov: } \vec{B} = \frac{\mu_0}{4\pi} \frac{q \vec{v} \times \hat{\mathbf{r}}}{r^2} \quad d\vec{B} = \frac{\mu_0}{4\pi} \frac{Id\vec{s} \times \hat{\mathbf{r}}}{r^2}$$

$$H\text{-felt rundt } \infty \text{ lang leder: } H_\theta = \frac{I}{2\pi r} \quad H\text{-felt i lang, tynn solenoide: } H = I \cdot n = I \cdot \frac{N}{\ell}$$

$$\text{Ohms lov: } V = RI, \quad R = \rho \frac{\ell}{A} = \frac{1}{\sigma} \frac{\ell}{A}; \quad P = VI$$

$\sigma \vec{E} = \vec{J}$ , der strømtetthet  $= \vec{J} = nq\vec{v}_d$  og  $\vec{v}_d = \mu \vec{E}$  = driftsfart.

$$\text{Induktans: } \mathcal{E} = -L \frac{dI}{dt} \quad \mathcal{E}_2 = -M_{21} \frac{dI_1}{dt}, \quad M_{21} = M_{12} \quad \text{Spoler: } L = N \frac{\Phi_B}{I} \quad U = \frac{1}{2} LI^2$$

Lenz lov: En indusert strøm er alltid slik at den forsøker å motvirke forandringen i den magnetiske fluks som er årsak til strømmen.

### Nablaoperatoren:

Kartesiske koordinater  $(x, y, z)$ , med enhetsvektorer henholdsvis  $\hat{i}, \hat{j}$  og  $\hat{k}$ :

$$\begin{aligned} \text{grad}V &= \vec{\nabla}V = \hat{i} \frac{\partial V}{\partial x} + \hat{j} \frac{\partial V}{\partial y} + \hat{k} \frac{\partial V}{\partial z} \\ \text{div} \vec{D} &= \vec{\nabla} \cdot \vec{D} = \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z} \\ \vec{\nabla}^2 V &= \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} \\ \text{curl} \vec{D} &= \vec{\nabla} \times \vec{D} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ D_x & D_y & D_z \end{vmatrix} \end{aligned}$$

Sylinderkoordinater  $(r, \phi, z)$ , med enhetsvektorer henholdsvis  $\hat{r}, \hat{\phi}$  og  $\hat{k}$ :

$$\begin{aligned} \vec{\nabla}V &= \hat{r} \frac{\partial V}{\partial r} + \hat{\phi} \frac{1}{r} \frac{\partial V}{\partial \phi} + \hat{k} \frac{\partial V}{\partial z} \\ \vec{\nabla} \cdot \vec{D} &= \frac{1}{r} \frac{\partial}{\partial r} (r D_r) + \frac{1}{r} \frac{\partial D_\phi}{\partial \phi} + \frac{\partial D_z}{\partial z} \\ \vec{\nabla}^2 V &= \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} \end{aligned}$$

Kulekoordinater  $(r, \theta, \phi)$ , med enhetsvektorer henholdsvis  $\hat{r}, \hat{\theta}, \hat{\phi}$ :

$$\begin{aligned} \vec{\nabla}V &= \hat{r} \frac{\partial V}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial V}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \\ \vec{\nabla} \cdot \vec{D} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 D_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (D_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial D_\phi}{\partial \phi} \\ \vec{\nabla}^2 V &= \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} \end{aligned}$$

Divergensteoremet og Stokes' teorem for et tilfeldig vektorfelt  $\vec{F}$ :

$$\begin{aligned} \iint \vec{F} \cdot d\vec{A} &= \iiint \vec{\nabla} \cdot \vec{F} \, d\tau \\ \oint \vec{F} \cdot d\vec{s} &= \iint (\vec{\nabla} \times \vec{F}) \cdot d\vec{A} \end{aligned}$$

Infinitesimale volumelement:

$$\begin{aligned} d\tau &= dx \, dy \, dz \\ d\tau &= r^2 dr \, \sin \theta \, d\theta \, d\phi \xrightarrow{\text{kulesymmetri}} 4\pi r^2 dr \\ d\tau &= r \, dr \, d\phi \, dz \xrightarrow{\text{syl.symmetri}} 2\pi r \, dr \, \ell \end{aligned}$$