Formelliste
TFY4155/FY1003 Elektrisitet og magnetisme

Formlenes gyldighetsområde og de ulike symbolenes betydning antas å være kjent. Symbolbruk som i forelesningene.
Siste revisjon: 21.4.16 ($|E| = c|B|$ justert fra versjon 13.4.15) (2 sider).

Q, ρ og σ uten indeks viser til *frie* ladninger. Q_i, ρ_i og σ_i er indusert ladning.
I og $\vec{J}$ uten indeks er ledningsstrøm (conducting current), I_d og $\vec{J}_\mathrm{d}$ er forskyvningsstrøm (displacement current).

Coulombs lov: $\vec{F}_{12} = \frac{1}{4\pi\epsilon}\frac{q_1 q_2}{r^2}\,\hat{\mathbf{r}}_{12}$ $\quad \vec{E} = \frac{1}{4\pi\epsilon}\frac{q}{r^2}\,\hat{\mathbf{r}}$

Gauss' lov integralform: $\oiint \vec{D}\cdot \mathrm{d}\vec{A} = Q$ $\quad \oiint \vec{E}\cdot \mathrm{d}\vec{A} = Q/\epsilon$ $\quad \oiint \vec{P}\cdot \mathrm{d}\vec{A} = -Q_\mathrm{i}$ $\quad \oiint \vec{B}\cdot \mathrm{d}\vec{A} = 0$

Gauss' lov differensialform: $\mathrm{div}\vec{D} = \rho$ $\quad \mathrm{div}\vec{E} = \rho/\epsilon$ $\quad \mathrm{div}\,\vec{P} = -\rho_\mathrm{i}$ $\quad \mathrm{div}\vec{B} = 0$

Fluks: $\Phi_E = \iint \vec{E}\cdot \mathrm{d}\vec{A}$ $\quad \Phi = \iint \vec{D}\cdot \mathrm{d}\vec{A} = \epsilon\,\Phi_E$ $\quad \Phi_B = \iint \vec{B}\cdot \mathrm{d}\vec{A}$

Amperes lov: $\oint \vec{B}\cdot \mathrm{d}\vec{s} = \mu\left(I + \epsilon\frac{\partial \Phi_E}{\partial t}\right)$ $\quad \oint \vec{H}\cdot \mathrm{d}\vec{s} = I + \frac{\partial \Phi}{\partial t}$ $\quad \mathrm{curl}\,\vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$

Faradays lov: $\mathcal{E} = -\frac{\partial \Phi_B}{\partial t} = -L\frac{\mathrm{d}I}{\mathrm{d}t}$ $\quad \oint \vec{E}\cdot \mathrm{d}\vec{s} = -\frac{\partial \Phi_B}{\partial t}$ $\quad \mathrm{curl}\,\vec{E} = -\frac{\partial \vec{B}}{\partial t}$

Maxwells likninger: $\mathrm{div}\vec{D} = \rho$ $\quad \mathrm{div}\vec{B} = 0$ $\quad \mathrm{curl}\,\vec{E} = -\frac{\partial \vec{B}}{\partial t}$ $\quad \mathrm{curl}\,\vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$

Forskyvningsstrøm: $I_\mathrm{d} = \frac{\partial \Phi}{\partial t}\,,\quad \vec{J}_\mathrm{d} = \frac{\partial \vec{D}}{\partial t}$

Elektrisk dipolmoment: $\vec{p} = q\vec{d}$ (fra − til +) $\quad$ Polarisering: $\vec{P} = \frac{\sum \vec{p}}{\text{volum}}$

Magnetisk (dipol)moment: $\vec{\mu} = \vec{m} = I\vec{A}$ $\quad$ Magnetisering: $\vec{M} = \frac{\sum \vec{\mu}}{\text{volum}}$

Kraftmoment: $\vec{\tau} = \vec{p}\times\vec{E}$ $\quad \vec{\tau} = \vec{\mu}\times\vec{B}$

$\vec{D} = \epsilon_0\vec{E} + \vec{P} = \epsilon\vec{E} = \epsilon_\mathrm{r}\epsilon_0\vec{E}$ $\quad \vec{P} = \chi_\mathrm{e}\epsilon_0\vec{E}$ $\quad \epsilon_\mathrm{r} = 1 + \chi_\mathrm{e}$

$\vec{B} = \mu_0\vec{H} + \mu_0\vec{M} = \mu\vec{H} = \mu_\mathrm{r}\mu_0\vec{H}$ $\quad \vec{M} = \chi_\mathrm{m}\vec{H}$ $\quad \mu_\mathrm{r} = 1 + \chi_\mathrm{m}$

Elektrisk potensial: $V_a - V_b = -\int_b^a \vec{E}\cdot \mathrm{d}\vec{s}\,,$ $\quad \vec{E} = -\vec{\nabla}V\,,$ $\quad$ Relativt ∞: $V(r) = \int \frac{\mathrm{d}q}{4\pi\epsilon r}$

Energi og energitetthet: $U = \frac{1}{2}\iiint V\,\mathrm{d}q$ $\quad$ Elektrisk: $u = \frac{1}{2}\vec{D}\cdot\vec{E}$ $\quad$ Magnetisk: $u = \frac{1}{2}\vec{B}\cdot\vec{H}$

Kondensatorer: $C = \frac{Q}{V}$ $\quad$ Kulekondensator: $C = 4\pi\epsilon_0 R$ $\quad$ Energi: $U = \frac{1}{2}QV = \frac{1}{2}CV^2$

Platekondensator: $C = \epsilon\frac{A}{d}$ $\quad$ Parallellkopling: $C = \sum_i C_i$ $\quad$ Seriekopling: $\frac{1}{C} = \sum_i \frac{1}{C_i}$

Kraft på strømførende leder: $\mathrm{d}\vec{F} = I\mathrm{d}\vec{s}\times\vec{B}$ $\quad$ Lorentzkrafta: $\vec{F} = q\left(\vec{E} + \vec{v}\times\vec{B}\right)$

Biot-Savarts lov: $\vec{B} = \frac{\mu_0}{4\pi}\frac{q\,\vec{v}\times\hat{\mathbf{r}}}{r^2}$ $\quad \mathrm{d}\vec{B} = \frac{\mu_0}{4\pi}\frac{I\mathrm{d}\vec{s}\times\hat{\mathbf{r}}}{r^2}$

H-felt rundt ∞ lang leder: $H_\theta = \frac{I}{2\pi r}$ $\quad$ H-felt i lang, tynn solenoide: $H = I\cdot n = I\cdot\frac{N}{\ell}$

Ohms lov: $V = RI\,,\quad R = \rho\frac{\ell}{A} = \frac{1}{\sigma}\frac{\ell}{A}\,;\quad P = VI$

$\sigma\vec{E} = \vec{J}\,,$ der strømtetthet $= \vec{J} = nq\vec{v}_\mathrm{d}$ og $\vec{v}_\mathrm{d} = \mu\vec{E}$ = driftsfart.

Induktans: $\mathcal{E} = -L\frac{\mathrm{d}I}{\mathrm{d}t}$ $\mathcal{E}_2 = -M_{21}\frac{\mathrm{d}I_1}{\mathrm{d}t}, \quad M_{21} = M_{12}$ Spoler: $L = N\frac{\Phi_B}{I}$ $U = \frac{1}{2}LI^2$

Lenz lov: En indusert strøm er alltid slik at den forsøker å motvirke forandringen i den magnetiske fluks som er årsak til strømmen.

Kompleks AC-signal: $V(t) = V_0 e^{i\omega t} = |V_0|e^{i\alpha}e^{i\omega t}$ $I(t) = I_0 e^{i\omega t} = |I_0|e^{i\beta}e^{i\omega t}$

$$Z = \frac{V(t)}{I(t)} = \frac{V_0}{I_0} = |Z|e^{i\phi} \qquad Z_R = R \qquad Z_L = i\omega L \qquad Z_C = \frac{1}{i\omega C}$$

Elektromagnetiske bølger:

Bølgelikningen for $\vec{E}$: $\nabla^2\vec{E} = \frac{1}{c^2}\frac{\partial^2\vec{E}}{\partial t^2}$ der $\nabla^2\vec{E} = \frac{\partial^2 E_x}{\partial x^2}\hat{\mathbf{i}} + \frac{\partial^2 E_y}{\partial y^2}\hat{\mathbf{j}} + \frac{\partial^2 E_z}{\partial z^2}\hat{\mathbf{k}}$ og $\frac{1}{c^2} = \mu\epsilon$

Bølge i $\pm x$-retning med $\vec{E}$ planpolarisert i y-retning: $\vec{E}(x,t) = E_0\,\hat{\mathbf{j}}\cos(\omega t \mp kx)$, $\vec{B}(x,t) = B_0\,\hat{\mathbf{k}}\cos(\omega t \mp kx)$

$\omega = 2\pi f$ $k = \frac{2\pi}{\lambda}$ $|c| = \frac{\omega}{k} = \sqrt{\frac{1}{\mu\epsilon}}$ $|E_0| = c|B_0|$ Bølge(vandre)retning som $\vec{E}\times\vec{B}$

Nablaoperatoren:

Kartesiske koordinater (x,y,z), med enhetsvektorer henholdsvis $\hat{\mathbf{i}}$, $\hat{\mathbf{j}}$ og $\hat{\mathbf{k}}$:

$$\begin{aligned}
\mathrm{grad}V = \vec{\nabla}V &= \hat{\mathbf{i}}\frac{\partial V}{\partial x} + \hat{\mathbf{j}}\frac{\partial V}{\partial y} + \hat{\mathbf{k}}\frac{\partial V}{\partial z}\\
\mathrm{div}\vec{D} = \vec{\nabla}\cdot\vec{D} &= \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z}\\
\vec{\nabla}^2 V &= \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}\\
\mathrm{curl}\,\vec{D} = \vec{\nabla}\times\vec{D} &= \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ D_x & D_y & D_z \end{vmatrix}
\end{aligned}$$

Sylinderkoordinater (r,ϕ,z), med enhetsvektorer henholdsvis $\hat{\mathbf{r}}$, $\hat{\phi}$ og $\hat{\mathbf{k}}$:

$$\begin{aligned}
\vec{\nabla}V &= \hat{\mathbf{r}}\frac{\partial V}{\partial r} + \hat{\phi}\frac{1}{r}\frac{\partial V}{\partial \phi} + \hat{\mathbf{k}}\frac{\partial V}{\partial z}\\
\vec{\nabla}\cdot\vec{D} &= \frac{1}{r}\frac{\partial}{\partial r}(rD_r) + \frac{1}{r}\frac{\partial D_\phi}{\partial \phi} + \frac{\partial D_z}{\partial z}\\
\vec{\nabla}^2 V &= \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial V}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2}
\end{aligned}$$

Kulekoordinater (r,θ,ϕ), med enhetsvektorer henholdsvis $\hat{\mathbf{r}}$, $\hat{\theta}$, $\hat{\phi}$:

$$\begin{aligned}
\vec{\nabla}V &= \hat{\mathbf{r}}\frac{\partial V}{\partial r} + \hat{\theta}\frac{1}{r}\frac{\partial V}{\partial \theta} + \hat{\phi}\frac{1}{r\sin\theta}\frac{\partial V}{\partial \phi}\\
\vec{\nabla}\cdot\vec{D} &= \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2 D_r\right) + \frac{1}{r\sin\theta}\frac{\partial}{\partial \theta}(D_\theta\sin\theta) + \frac{1}{r\sin\theta}\frac{\partial D_\phi}{\partial \phi}\\
\vec{\nabla}^2 V &= \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial V}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial \theta}\left(\sin\theta\frac{\partial V}{\partial \theta}\right) + \frac{1}{r^2\sin^2\theta}\frac{\partial^2 V}{\partial \phi^2}
\end{aligned}$$

Divergensteoremet og Stokes' teorem for et tilfeldig vektorfelt $\vec{F}$:

$$\oiint \vec{F}\cdot\mathrm{d}\vec{A} = \iiint \vec{\nabla}\cdot\vec{F}\;\mathrm{d}\tau \qquad \oint \vec{F}\cdot\mathrm{d}\vec{s} = \iint \left(\vec{\nabla}\times\vec{F}\right)\cdot\mathrm{d}\vec{A}$$

Infinitesimale volumelement:

$$\begin{aligned}
\mathrm{d}\tau &= \mathrm{d}x\,\mathrm{d}y\,\mathrm{d}z\\
\mathrm{d}\tau &= r^2\,\mathrm{d}r\,\sin\theta\,\mathrm{d}\theta\,\mathrm{d}\phi \xrightarrow{\text{kulesymmetri}} 4\pi r^2\,\mathrm{d}r\\
\mathrm{d}\tau &= r\,\mathrm{d}r\,\mathrm{d}\phi\,\mathrm{d}z \xrightarrow{\text{syl.symmetri}} 2\pi r\,\mathrm{d}r\,\ell
\end{aligned}$$