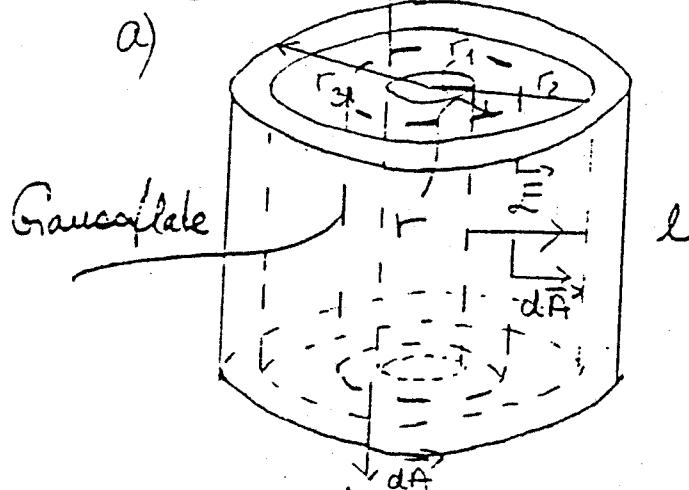


ØSNINGSFORSLAG EKSAHEN AUG 1997

FAS 74142 FYSIKK

Oppgave 1



\vec{E} radielt rettet ut fra metallbaren (der bort fra endeffekter i hver ende av kondensatoren)

- i) Gaußflaten velges som en sylinder med endeffekter i området der feltet skal beregnes (Snitt i figuren over)
 $\vec{E} \parallel d\vec{A}$ på sylinderflaten og $\vec{E} \perp d\vec{A}$ på endeffektene.

Gauss lov for $r_1 < r < r_2$:

$$\oint \vec{E} \cdot d\vec{A} = \int_{\text{syl.fl.}} \vec{E} \cdot d\vec{A} + 2 \int_{\text{endell.}} \vec{E} \cdot d\vec{A} = \int_{\text{syl.fl.}} \vec{E} \cdot d\vec{A}$$

$\underbrace{\vec{E} \perp d\vec{A} \Rightarrow 0}_{\vec{E} \parallel d\vec{A}}$ og $|\vec{E}| = \text{konst}$
 på syl. fl.
 pga. symmetri

$$\Rightarrow \oint \vec{E} \cdot d\vec{A} = E \int_{\text{syl.fl.}} dA = E \cdot 2\pi r l = \frac{Q_{\text{innenfor}}}{\epsilon_0}$$

$$Q_{\text{innenfor}} = Q \Rightarrow E = \frac{Q}{2\pi\epsilon_0 r l} \quad \text{q.e.d}$$

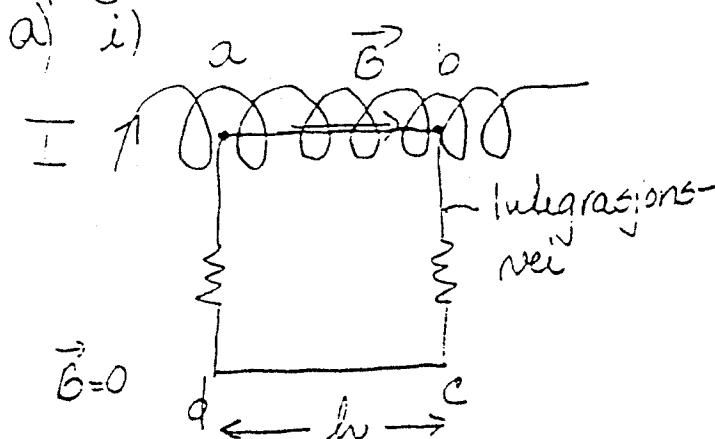
- ii) Inn i metall: $E = 0$

$$\Rightarrow E = 0 \text{ for } r < r_1 \text{ og } r_2 < r < r_3$$

$$\text{For } r > r_3 : Q_{\text{innenfor}} = Q - Q = 0$$

$$\Rightarrow E = 0 \text{ for } r > r_3$$

Oppgave 2



Ampères lov:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{innerfor}}$$

Legger en lukket integrasjonsvei som viser på figuren. Linjeslykket cd er lang unna spolen slik at magnetfeltet her er $\vec{B} = 0$.

$$\Rightarrow \oint \vec{B} \cdot d\vec{l} = \int_{ab} \vec{B} \cdot d\vec{l} + \int_{bc} \vec{B} \cdot d\vec{l} + \int_{cd} \vec{B} \cdot d\vec{l} + \int_{da} \vec{B} \cdot d\vec{l}$$

$$= \int_{ab} B dl = B \cdot h = \mu_0 I_{\text{innerfor}} = \mu_0 \underbrace{\frac{I \cdot N}{l}}_{I \cdot \text{antall vindinger pr. lengdeenhet}} \cdot h$$

$$\Rightarrow B = \frac{\mu_0 N I}{l} \quad \text{q.e.d} \quad \text{Feltretn. visst i figuren over.}$$

ii) Spoles selvinduktans: $L = \frac{N \Phi_B}{I}$

$$\text{der } \Phi_B = \int \vec{B} \cdot d\vec{A} = B \cdot A = B \cdot \frac{\pi d^2}{4}$$

B = konst
inni spoleversnitt

$$\Rightarrow L = \frac{N \mu_0 N I}{l} \cdot \frac{\pi d^2}{4} = \frac{\mu_0 N^2 \pi d^2}{4l} \quad (\text{Utklart ved dimensjoner, antall vindinger og } \mu_0)$$

b) Diff. lign.: $\frac{d^2 q}{dt^2} + \frac{1}{LC} q = 0$

med løsn. på formen $q = q_0 \cos(\omega t + \phi)$

$$i) \frac{dq}{dt} = -\frac{q_0}{T_0} \omega \sin(\omega t + \phi)$$

$$\frac{d^2q}{dt^2} = -\frac{q_0}{T_0} \omega^2 \cos(\omega t + \phi)$$

$$\text{Insatte i diff. lign.: } -\frac{q_0}{T_0} \omega^2 \cos(\omega t + \phi) + \frac{1}{LC} \frac{q_0}{T_0} \cos(\omega t + \phi) = 0$$

$$\Rightarrow \omega^2 = \frac{1}{LC} \quad \Rightarrow \underline{\underline{\omega = \frac{1}{\sqrt{LC}} \text{ q.e.d}}}$$

$$ii) \text{ Elektrisk feltenergi: } U_E = \frac{1}{2} C V^2 = \frac{1}{2} C q^2 = \frac{1}{2C} q_0^2 \cos^2(\omega t + \phi)$$

$$\text{Magnetisk feltenergi: } U_B = \frac{1}{2} L I^2 = \frac{1}{2} L \left(\frac{dq}{dt}\right)^2 = \frac{1}{2L} q_0^2 \omega^2 \sin^2(\omega t + \phi)$$

$$\text{Total feltenergi: } U = U_E + U_B$$

$$U = \frac{1}{2C} q_0^2 \cos^2(\omega t + \phi) + \frac{1}{2L} q_0^2 \omega^2 \sin^2(\omega t + \phi)$$

$$\omega^2 = \frac{1}{LC} \Rightarrow \underline{\underline{U = \frac{1}{2C} q_0^2 [\cos^2(\omega t + \phi) + \sin^2(\omega t + \phi)] = \frac{1}{2C} q_0^2 = \text{konst.}}} \text{ q.e.}$$

$$c) i) \text{ Startbedingelser: } \textcircled{I} I(t=0) = \frac{dq}{dt}_{t=0} = 0$$

$$\textcircled{II} q(t=0) = Q; \text{ ladu. på kondensator}$$

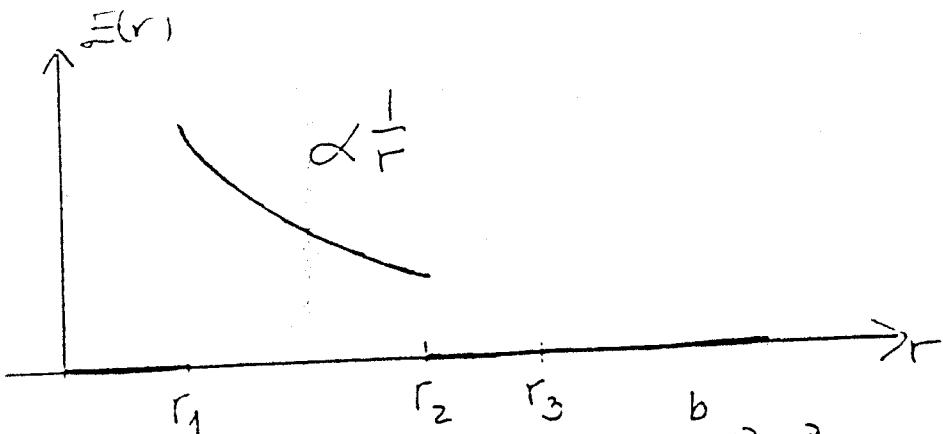
$$\textcircled{I} \Rightarrow -q_0 \omega \sin \phi = 0$$

$$q_0 \omega \neq 0 \Rightarrow \sin \phi = 0 \Rightarrow \underline{\underline{\phi = 0^\circ}}$$

$$\textcircled{II} \Rightarrow q_0 \cos \phi = \underline{\underline{q_0 = Q}}$$

(2)

iii)



b) i) Elektrisk potensial $V_b - V_a = - \int_a^b \vec{E} \cdot d\vec{l}$

$$r > r_2 \quad V(r) - V(r_2) = - \int_{r_2}^r \vec{E} \cdot d\vec{l} = 0 \quad \text{da } \vec{E} = 0$$

Velg referansepunkt $V(r_2) = 0 \Rightarrow V(r) = 0 \text{ for } r \geq r_2$

$$r_1 < r < r_2 \quad V(r_2) - V(r_1) = - \int_{r_1}^{r_2} \vec{E} \cdot d\vec{l} = - \int_{r_1}^{r_2} \vec{E} dt \quad \text{da } \vec{E} \parallel \vec{dl} = d\vec{r}$$

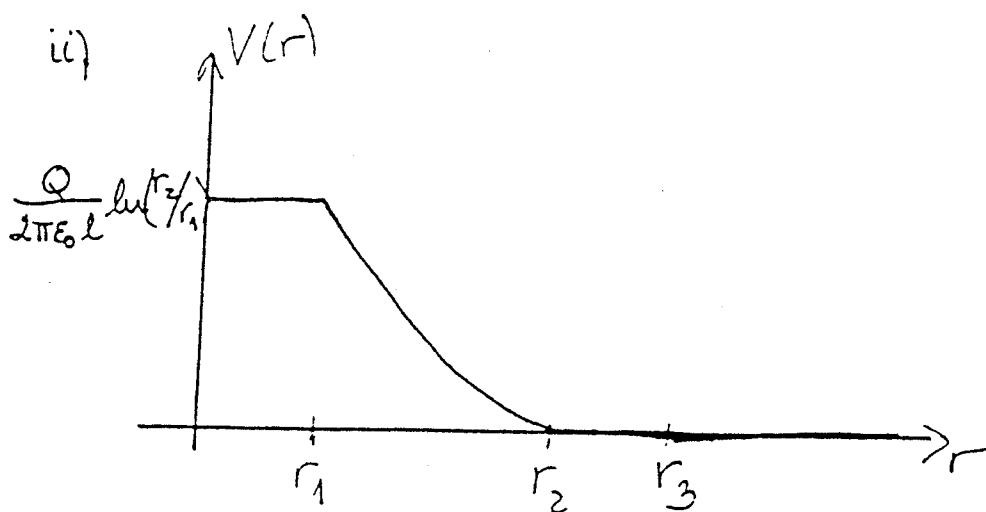
$$-V(r) = - \int_r^{r_2} \frac{Q}{2\pi\epsilon_0 r l} dr' = - \left[\frac{Q}{2\pi\epsilon_0 l} \ln r' \right]_r^{r_2}$$

$$\underline{V(r) = \frac{Q}{2\pi\epsilon_0 l} \ln(r_2/r)} \quad r_1 \leq r \leq r_2$$

$$r < r_1 \quad V(r_1) - V(r) = - \int_r^{r_1} \vec{E} \cdot d\vec{l} = 0 \quad \text{da } \vec{E} = 0$$

$$\Rightarrow \underline{V(r) = V(r_1) = \frac{Q}{2\pi\epsilon_0 l} \ln(r_2/r_1)} \quad r \leq r_1$$

ii)



c) i) ~~Kao~~ ~~taus~~ $C = \frac{Q}{\Delta V}$

ΔV er potensialforskjellen mellom metallsparen og sylinderen

$$\Delta V = V(r_1) - V(r_2) = \frac{Q}{2\pi\epsilon_0 l} \ln\left(\frac{r_2}{r_1}\right)$$

$$\Rightarrow \underline{\underline{C}} = \frac{\underline{\underline{Q}}}{\frac{\underline{\underline{Q}}}{2\pi\epsilon_0 l} \ln\left(\frac{r_2}{r_1}\right)} = \frac{\underline{\underline{\frac{Q}{2\pi\epsilon_0 l}}}}{\underline{\underline{\ln\left(\frac{r_2}{r_1}\right)}}} = \underline{\underline{0,56 \cdot 10^{-9} F}}$$

ii) Energiinnholdet i kondensatoren:

$$\underline{\underline{U}} = \frac{1}{2} C V^2 = \frac{1}{2} C \Delta V^2 = \frac{\underline{\underline{\frac{Q^2}{4\pi\epsilon_0 l}}}}{\underline{\underline{\ln\left(\frac{r_2}{r_1}\right)}}} = \underline{\underline{1,12 \cdot 10^{-3} J}}$$

d)

Vi kan se på systemet som en parallekkobling av to kondensatorer, $C = C_1 + C_2$

$$\Rightarrow C = \frac{2\pi\epsilon_0 l/2}{\ln r_2/r_1} + \frac{2\pi\epsilon_0 \epsilon_r l/2}{\ln r_2/r_1}$$

$$= \frac{2\pi\epsilon_0 l}{\ln r_2/r_1} \left(\frac{1}{2} + \frac{1}{2} \epsilon_r \right)$$

$$U = \frac{1}{2} C V^2 = U_c \left(\frac{1}{2} + \frac{1}{2} \epsilon_r \right)$$

Her er U_c energiinnholdet i oppgave c

$$\underline{\underline{U}} = \underline{\underline{1,12 \cdot 10^{-3} \left(\frac{1}{2} + \frac{1}{2} \cdot 80 \right)}} = \underline{\underline{45,92 \cdot 10^{-3} J}}$$

(ii) Vi har funnet $q_0 = Q$ $\phi = 0$

\Rightarrow Lösning:

$$q = Q_0 \cos \omega t$$

For $t = T/6$

$$q = Q_0 \cos 2\pi/6 = Q_0/2$$

$$\phi = 0$$

$$\text{Elektrisk felt energi } U_E = \frac{1}{2C} q_0^2 \cos^2 \omega t$$

$$\begin{aligned} \text{Magnetisk } &\sim - & U_B &= \frac{1}{2} \omega^2 L q_0^2 \sin^2 \omega t \\ &&&= \frac{1}{2C} q_0^2 \sin^2 \omega t \end{aligned}$$

$$\Rightarrow \underline{\underline{r}} = \frac{U_E}{U_B} = \frac{\cos^2 \omega t}{\sin^2 \omega t} = \frac{\left(\frac{1}{2}\right)^2}{\left(\frac{1}{2}\sqrt{3}\right)^2} = \underline{\underline{\frac{1}{3}}}$$

(6)

Opgave 3

a) i) Fenomenet kallas (Fraunhofer) diffraction

ii) Nullpunkt föreller i intensitetsfördelningen sit null:

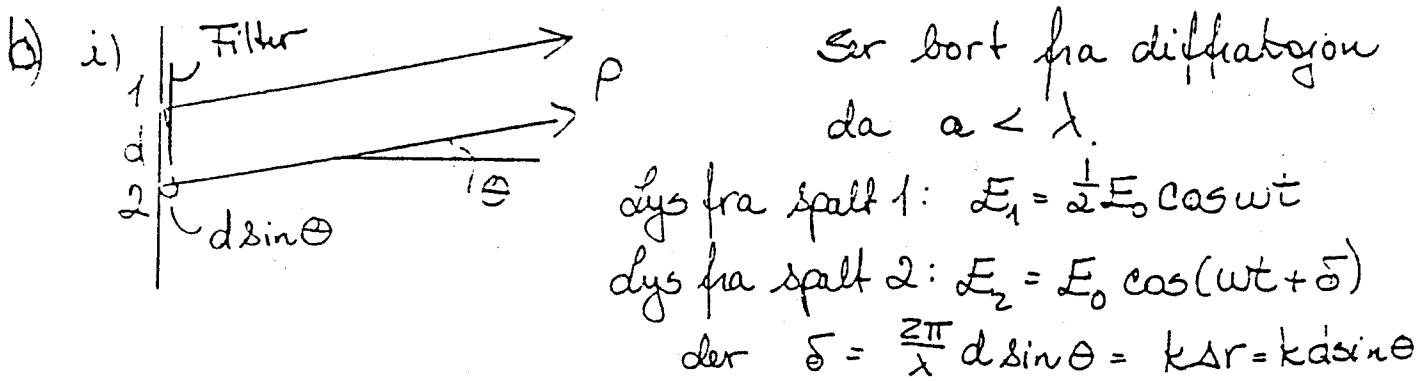
$$\sin(\beta/a) = 0$$

$$\Rightarrow \text{For } \beta/a = n\pi = \frac{\pi a}{\lambda} \sin\theta \text{ med } n = \pm 1, \pm 2, \dots$$

$$\Rightarrow \underline{\text{Nullpunkt för } \sin\theta = \frac{n\lambda}{a} \quad n = \pm 1, \pm 2, \dots}$$

(Inte nullpkt. för $n=0$ då $\frac{\sin \beta/a}{\beta/a} \rightarrow 1$ när $\beta/a \rightarrow 0$,
som leder till särskilsmats. β/a i diff. monstret)

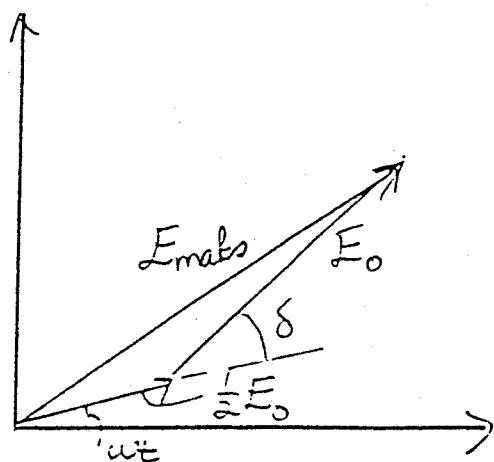
iii. När spaltebredden a ökar blir avståndet mellan
nullptt. mindre da $\sin\theta \propto \frac{1}{a}$



Superpositionsprincipet för lyset i ptt. P:

$$E = E_1 + E_2$$

$$E = \frac{1}{2} E_0 \cos \omega t + E_0 \cos(\omega t + \delta)$$



Bruter viserdiagram för
å finna resulterende
amplitud:

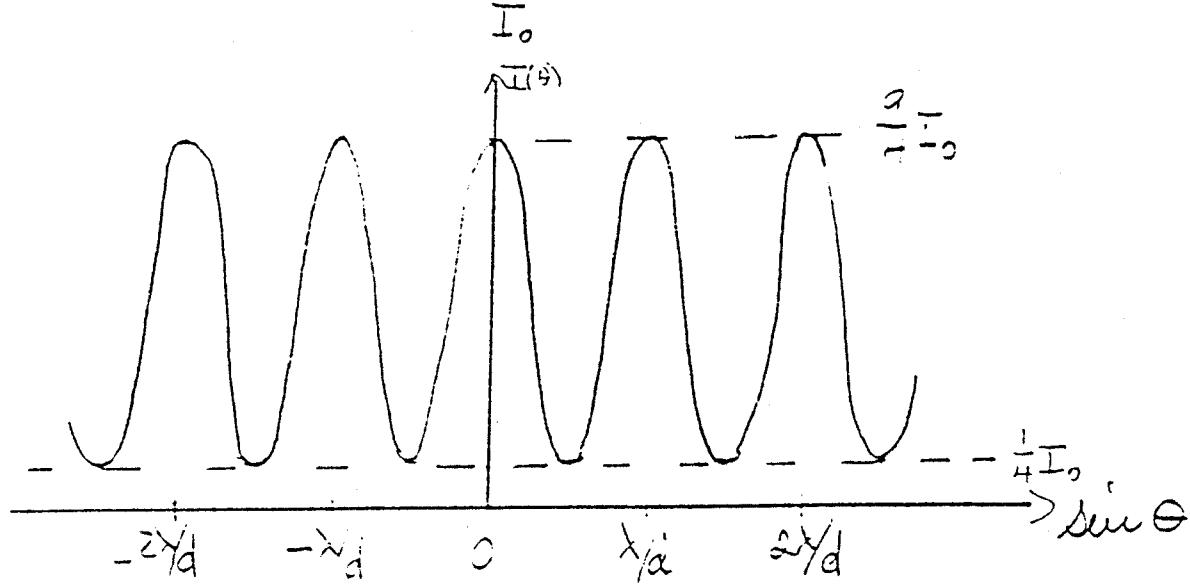
$$E_{\text{maks}}^2 = \frac{1}{2} E_0^2 + E_0^2 - 2 \cdot \frac{1}{2} E_0 E_0 \cos(\pi -$$

$$E_{\text{maks}}^2 = \frac{5}{4} E_0^2 + E_0^2 \cos \delta$$

$$\text{Intensitet} \quad I = \frac{1}{2} C \epsilon_0 \epsilon_m^2 = \frac{1}{2} C \epsilon_0 \epsilon_0^2 \left(\frac{5}{4} \cos \delta \right) \quad (8)$$

$$I = \underbrace{\frac{1}{2} C \epsilon_0 \epsilon_0^2}_{I_0} \left(\frac{5}{4} + \cos \left[\frac{2\pi}{\lambda} d \sin \theta \right] \right)$$

ii)



Sammenligning med dobbeltspalt:

Intensitet i hovedaksen $\frac{9}{4} I_0$ mot $4 I_0$ for dobbeltspalt

Intensitet i minima $\frac{1}{4} I_0$ mot 0 for dobbeltspalt.

Ellers litt. Forstjeller skyldes ulike amplituder på de to bølgende når filteret er til side.