

①

Eksamen, MNFFY-~~43~~/V99. Hvite dverger, 4/6-99  
 (MNFFY-450. Kompakte stjerner)

Antydning. Oppgave 1

a) Ved likevekt gjelder

$$\mu_e + \mu_p = \mu_n,$$

for de kjemiske potensialene, når vi setter

$$\mu_\nu \approx 0.$$

Med  $\mu = \epsilon_F = [(p_F c)^2 + (m c^2)^2]^{1/2} = m c^2 [1 + (p_F/mc)^2]^{1/2}$ ,  
 $x = p_F/mc$ ,

for henholdsvis elektron, proton og neutron, får vi

$$m_e (1+x_e^2)^{1/2} + m_p (1+x_p^2)^{1/2} = m_n (1+x_n^2)^{1/2}.$$

Ladningsneutralitet krever

$$n_e = x_e^3 / 3\pi^2 \lambda_e^3 = n_p = x_p^3 / 3\pi^2 \lambda_p^3,$$

$$m_e x_e = m_p x_p,$$

der  $\hbar/mc = \lambda$

er Compton-bølglengden. Blandingsforholdet mellom partiklene finnes nå ved innsettning i ligningen for de kjemiske potensialene, dvs.

$$\underline{(m_e^2 + m_p^2 x_p^2)^{1/2} + m_p (1+x_p^2)^{1/2} = m_n (1+x_n^2)^{1/2}}$$

b) Kvadrering av det siste uttrykket under a) gir de

$$m_e^2 + m_p^2 x_p^2 = [m_n (1+x_n^2)^{1/2} - m_p (1+x_p^2)^{1/2}]^2$$

$$= m_n^2 + m_n^2 x_n^2 - 2m_n m_p (1+x_n^2)^{1/2} (1+x_p^2)^{1/2} + m_p^2 x_p^2 + m_p^2,$$

$$(m_n^2 + m_p^2 - m_e^2 + m_n^2 x_n^2)^2 = 4m_n^2 m_p^2 (1+x_n^2)(1+x_p^2)$$

$$= 4m_n^2 m_p^2 (1+x_n^2 + x_p^2 + x_n^2 x_p^2),$$

$$(m_n^2 + m_p^2 - m_e^2)^2 + 2m_n^2 x_n^2 (m_n^2 + m_p^2 - m_e^2) + m_n^4 x_n^4$$

$$= 4m_n^2 m_p^2 x_p^2 (1+x_n^2) + 4m_n^2 m_p^2 (1+x_n^2),$$

$$4 m_n^2 m_p^2 x_p^2 (1 + x_n^2) = (Q^2 - m_e^2) [(m_n + m_p)^2 - m_e^2] + 2 m_n^2 x_n^2 (m_n^2 - m_p^2 - m_e^2) + m_n^4 x_n^4$$

sidan  $Q^2 - m_e^2 = (m_n - m_p)^2 - m_e^2$ ,

dvs. 
$$n_p/n_n = \frac{(m_p x_p / m_n x_n)^3}{\left[ \left( 1 + 2 \frac{(m_n^2 - m_p^2 - m_e^2)}{m_n^2 x_n^2} + \frac{(Q^2 - m_e^2)}{m_n^2 x_n^2} \right) \times \frac{[(m_n + m_p)^2 - m_e^2]}{m_n^4 x_n^4} \right]^{3/2} / (1 + 1/x_n^2)} / 8}$$

$$\approx \frac{\left[ 1 + 4Q/m_n x_n^2 + 4(Q^2 - m_e^2)/m_n^2 x_n^4 \right]^{3/2} / 8}{(1 + 1/x_n^2)}$$

sidan  $Q = m_n - m_p \ll m_p \approx m_n$   
 $m_e \ll m_p \approx m_n$   
 $m_n^2 - m_p^2 - m_e^2 = (m_n - m_p)(m_n + m_p) - m_e^2 \approx 2Q m_n$

c) Vi får minimum for  $\partial(n_p/n_n)/\partial x_n = 0$ ,  $(n_n = x_n^3 / 3\pi^2 \lambda_n^3)$

dvs. for  $\partial \left\{ \frac{[x_n^4 + 4Qx_n^2/m_n + 4(Q^2 - m_e^2)/m_n^2]}{(x_n^4 + x_n^2)} \right\} / \partial x_n = 0$ ,  
 $(4x_n^3 + 8Qx_n/m_n)(x_n^4 + x_n^2) - [x_n^4 + 4Qx_n^2/m_n + 4(Q^2 - m_e^2)/m_n^2] \times (4x_n^3 + 2x_n) = 0$

$$4x_n^6 + 8Qx_n^4/m_n + 4x_n^4 + 8Qx_n^2/m_n - 4x_n^6 - 16Qx_n^4/m_n - 16(Q^2 - m_e^2)x_n^2/m_n^2 - 2x_n^4 - 8Qx_n^2/m_n - 8(Q^2 - m_e^2)/m_n^2 = 0$$

$$n_n \approx \frac{(2^{3/2} / 3\pi^2 \lambda_n^3) [(Q^2 - m_e^2)/m_n^2]^{3/4}}{2^{1/2} [(Q^2 - m_e^2)/m_n^2]^{1/4}}$$

som gir  $(n_p/n_n)_{\min} = \left\{ (Q/m_n) + [(Q^2 - m_e^2)^{1/2}/m_n] \right\}^{3/2} \left( \frac{Q \ll m_n}{x_n \ll 1} \right)$

For meget store tettheter vil da

$$x_n \rightarrow \infty, \text{ for } \rho_0 \rightarrow \infty,$$

$$(n_p/n_n) \rightarrow 1/8, \text{ for } x_n \rightarrow \infty.$$

Dessuten gjelder

$$n_e = n_p$$

ved ladningsneutralitet, dvs.

$$\underline{n_e : n_p : n_p \rightarrow 1 : 1 : 8, \text{ for } \rho_0 \rightarrow \infty}$$

## Antydret løsning. Oppgave 2

a) Ligningene

$$\begin{aligned} dm/dr &= 4\pi r^2 \rho, \\ dP/dr &= -6m\rho/r^2, \\ P &= K\rho^\Gamma, \end{aligned}$$

kan kombineres til

$$\begin{aligned} m &= -(r^2/\rho 6)(dP/dr), \\ dm/dr &= -(1/6)d[(r^2/\rho)/(dP/dr)]/dr = 4\pi r^2 \rho, \\ dP/dr &= \Gamma K \rho^{\Gamma-1} (d\rho/dr), \end{aligned}$$

dvs.  $r^{-2} d[(r^2/\rho)(dP/dr)]/dr = (\Gamma K/r^2) d[r^2 \rho^{\Gamma-2} (d\rho/dr)]/dr = -4\pi 6 \rho,$

$$\begin{aligned} r^2 d[r^2 \rho^{\Gamma-2} (d\rho/dr)]/dr &= -(4\pi 6/\Gamma K) \rho \\ &= a^{-2} \xi^{-2} d[a^2 \xi^2 \rho_c^{\Gamma-2} \theta^{(\Gamma-2)n} (d\rho/d\xi)]/d\xi \\ &= a^{-3} \xi^{-2} d[\rho_c a^2 \xi^2 \theta^{n(\Gamma-2)} (d\rho/d\xi)]/d\xi \\ &= (\rho_c^{\Gamma-2}/a^2 \xi^2) d[\xi^2 \theta^{n(\Gamma-2)} \rho_c n \theta^{n-1} (d\theta/d\xi)]/d\xi \\ &= (n \rho_c^{\Gamma-1}/a^2 \xi^2) d[\xi^2 \theta^{n(\Gamma-1)-1} (d\theta/d\xi)]/d\xi = -(4\pi 6/\Gamma K) \rho_c \theta^n, \\ \xi^{-2} d[\xi^2 \theta^{n(\Gamma-1)-1} (d\theta/d\xi)]/d\xi &= -(4\pi 6 a^2/\Gamma K n \rho_c^{\Gamma-2}) \theta^n. \end{aligned}$$

Tilpasning til Lane-Emden's ligning gir oss da betingelsene

$$\begin{aligned} n\Gamma - n - 1 &= 0, \\ 4\pi 6 a^2/\Gamma K n \rho_c^{\Gamma-2} &= 1, \end{aligned}$$

dvs.  $\Gamma = (n+1)/n = 1 + (1/n), \quad n = 1/(\Gamma-1),$

$$a = (\Gamma K n \rho_c^{\Gamma-2}/4\pi 6)^{1/2} = [(n+1)K \rho_c^{(n-1)/n}/4\pi 6]^{1/2},$$

$$\xi^{-2} d[\xi^2 (d\theta/d\xi)]/d\xi = -\theta^n.$$

b) Total masse er gitt ved

$$\begin{aligned} M &= \int_0^R 4\pi r^2 \rho dr = \int_0^{\xi_1} 4\pi a^2 \xi^2 \rho_c \theta^n a d\xi = 4\pi a^3 \rho_c \int_0^{\xi_1} \xi^2 \theta^n d\xi \\ &= -4\pi a^3 \rho_c \int_0^{\xi_1} [d(\xi^2 d\theta/d\xi)/d\xi] d\xi = 4\pi a^3 \rho_c \xi_1^2 |\theta'(\xi_1)|, \end{aligned}$$

dvs.

$$M = 4\pi [(n+1)K/4\pi G]^{3/2} \rho_c^{(3-n)/2n} \frac{e^2}{\xi_1} |\Theta'(\xi_1)|.$$

Total radius blir

$$R = a\xi_1 = [(n+1)K/4\pi G]^{1/2} \rho_c^{(1-n)/2n} \frac{e}{\xi_1}.$$

c) Fra uttrykket for total masse  $M$ , finner vi at forholdet mellom midlere tetthet  $\bar{\rho}$  og sentral tetthet  $\rho_c$  blir lik

$$\begin{aligned} \bar{\rho}/\rho_c &= (3M/4\pi R^3)/\rho_c \\ &= 3 \cdot 4\pi a^3 \rho_c \xi_1^2 |\Theta'(\xi_1)| / (4\pi a^3 \xi_1^3 \rho_c) = \underline{3|\Theta'(\xi_1)|/\xi_1}, \end{aligned}$$

siden

$$M = (4\pi R^3/3) \bar{\rho}.$$

Fra uttrykket for total radius  $R$ , finner vi at

$$\begin{aligned} R &\sim \rho_c^{(1-n)/2n} \frac{e}{\xi_1}, \\ \rho_c &\sim R^{2n/(1-n)} \end{aligned}$$

og fra uttrykket for total  $M$  får vi

$$M \sim a^3 \rho_c \frac{e^2}{\xi_1} \sim \rho_c^{(3-n)/2n} \sim R^{(3-n)/(1-n)}.$$
For  $\Gamma = 5/3$ 

får vi da

$$n = 1/(\Gamma - 1) = 3/2,$$

$$\underline{M \sim R^{-(3/2)/(1/2)} = R^{-3}}$$

### Antydning løsning, Oppgave 3

a) Poissons ligning gir

$$\nabla^2 \Phi = 4\pi G \rho,$$

dvs.  $\nabla^2 \delta \Phi = 4\pi G \delta \rho,$

der  $\delta \rho = -\nabla \cdot (\rho \xi),$

dvs.  $\delta \Phi = -G \int \frac{\delta \rho'}{|\underline{x} - \underline{x}'|} d^3 x' = G \int \frac{\nabla' \cdot (\rho' \xi')}{|\underline{x} - \underline{x}'|} d^3 x' = -G \int \rho' \xi' \cdot \nabla' \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) d^3 x',$

ifølge delvis integrasjon, siden

$$\rho' = 0 \text{ for } \underline{x}' \in \mathbb{R},$$

på overflaten. Dessuten er

$$\nabla \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) = -\nabla' \left( \frac{1}{|\underline{x} - \underline{x}'|} \right),$$

$$\nabla^2 \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) = -4\pi \delta(\underline{x} - \underline{x}').$$

Derivasjon av uttrykket for  $\delta \Phi$  gir da

$$\begin{aligned} \nabla \delta \Phi &= \nabla \left( G \int \frac{\nabla' \cdot (\rho' \xi')}{|\underline{x} - \underline{x}'|} d^3 x' \right) = G \int \nabla' (\rho' \xi') \cdot \nabla \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) d^3 x' \\ &= -G \int \nabla' (\rho' \xi') \cdot \nabla' \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) d^3 x' \\ &= -G \left[ \rho' \xi' \cdot \nabla' \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) + \int \rho' \xi' \cdot \nabla'^2 \left( \frac{1}{|\underline{x} - \underline{x}'|} \right) d^3 x' \right] \\ &= -4\pi G \int \rho' \xi' \cdot \delta(\underline{x} - \underline{x}') d^3 x' = -4\pi G \rho \xi_i. \end{aligned}$$

dvs.  $\nabla_i \delta \Phi = -4\pi G \rho \xi_i$

Siden  $\nabla_i \delta \Phi = (x_i/r) \nabla \delta \Phi,$   
 $\xi_i = (x_i/r) \xi,$

b) Ligninger for hydrostatisk likevekt følger egentlig fra ligninger for impulsbevarelse, som gir

$$\nabla p + \rho \nabla \Phi = 0,$$

og ligninger for radiale oscillasjoner på grunn av

somå perturbationer av en statisk likevekt for en kulesymmetrisk stjerne, blir da

$$\begin{aligned}
 \sum_j L_{ij} \xi_j &= \sum_j \nabla_i (\Gamma_i P \nabla_j \xi_j) - (\nabla_j \xi_j) \nabla_i P + (\nabla_i \xi_j) \nabla_j P \\
 &\quad - \rho \xi_j \nabla_j \nabla_i \Phi - \rho \nabla_i \delta \Phi \\
 &= d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\}/dr - [d(r^2 \xi_j)/dr] (dP/dr)/r^2 \\
 &\quad + (d\xi_j/dr) (dP/dr) - \rho \xi_j d^2 \Phi/dr^2 + 4\pi G \rho^2 \xi_j \\
 &= d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\}/dr - [(2/r) \xi_j + (d\xi_j/dr)] (dP/dr) \\
 &\quad + (d\xi_j/dr) (dP/dr) - \rho \xi_j [\nabla^2 \Phi - (2/r) (d\Phi/dr)] + 4\pi G \rho^2 \xi_j \\
 &= d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\}/dr - (2/r) (dP/dr) \xi_j - 4\pi G \rho^2 \xi_j \\
 &\quad - (2/r) (dP/dr) \xi_j + 4\pi G \rho^2 \xi_j \\
 &= d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\}/dr - (4/r) (dP/dr) \xi_j,
 \end{aligned}$$

som gir oss ligningen

$$\underline{d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\} - (4/r) (dP/dr) \xi_j + \omega^2 \rho \xi_j = 0.}$$

c) For  $\Gamma_i = 4/3$ ,

far vi egenverdi-ligningen

$$\begin{aligned}
 &d\{\Gamma_i P [d(r^2 \xi_j)/dr]/r^2\} - (4/r) (dP/dr) \xi_j + \omega^2 \rho \xi_j \\
 &= d\{(4P/3) [d(K \cdot r^3)/dr]/r^2\} - (4/r) (dP/dr) K r + \omega^2 \rho r \\
 &= d(4KP)/dr - 4K (dP/dr) + \omega^2 \rho r = \omega^2 \rho r = 0,
 \end{aligned}$$

dvs.

$$\underline{\omega^2 = 0,}$$

for

$$\underline{\xi_j = K \cdot r = \text{konstant} \cdot r.}$$